

Blow-up phenomena in a half-space

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Main goal

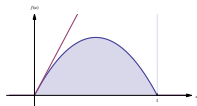
Determine if solutions, starting from $u_0 \geq 0$, to

$$\partial_t u = \mathcal{D}[u] + u^{1+p}, \quad t > 0, \mathbf{x} \in \Omega, \quad (p > 0),$$

are global in time or not.

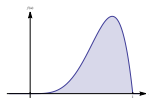
► Direct implications on the population dynamics models

$$\partial_t u = \mathcal{D}[u] + u^{1+p}(1-u),$$



KPP ($p = 0$)

vs.



vs. Allee effect ($p > 0$)

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- 2 In the half-space $\mathbb{R}_+^N = \mathbb{R}^{N-1} \times (0, +\infty)$
- 3 The road-field model

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The Heat equation

The solution $v(t, x)$ to

$$\partial_t v = \Delta v, \quad v(t=0, \cdot) = v_0,$$

is given by

$$v(t, x) = G(t, \cdot) * v_0(x).$$

It is global and satisfies

$$\|v(t, \cdot)\|_{L^\infty(\mathbb{R}^N)} \leq \frac{C(v_0)}{(1+t)^{N/2}}.$$

Nonlinear ODE: systematic blow up

For $p > 0$, the solution $u(t)$ to

$$\frac{du}{dt} = u^{1+p}, \quad u(0) = u_0 > 0,$$

always blows up in finite time.

$$\partial_t u = \Delta u + u^{1+p}, \quad p > 0$$

Starting from a nontrivial compactly supported $u_0 \geq 0$, what happens to the Cauchy problem?

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Theorem (Fujita 1966)

Define $p_F := \frac{2}{N}$.

- (i) $0 < p \leq p_F \implies$ all solutions blow up in finite time.
- (ii) $p > p_F \implies$ some solutions ("small" u_0) are global and get extinct "like the Heat equation".

Proof for the case $p > p_F$

► Look for a supersolution in the form $u^+(t, x) := g(t)v(t, x)$ where $v := e^{t\Delta}u_0$, $g(0) = 1$, and hope that g is global...

► This requires

$$\frac{g'(t)}{g^{1+p}(t)} \geq v^p(t, x).$$

► It is enough to select

$$\frac{g'(t)}{g^{1+p}(t)} = \left(\frac{C(u_0)}{(1+t)^{N/2}} \right)^p, \quad g(0) = 1,$$

whose solution is computable and global (thanks to $p > \frac{2}{N}$) if “ u_0 is small enough”.

Proof for the case $p < p_F$

Assume u is global and consider $f(t) := \int_{\mathbb{R}^N} G(t, y) u_0(y) dy$.

► **Linear argument:** $f(t)$ is nothing else than $v(t, 0)$ where $v = e^{t\Delta} u_0$ so that

$$f(t) \gtrsim \frac{1}{t^{N/2}}.$$

► **Nonlinear argument:** $f(t)$ is nothing else than $g(0)$ where

$$g(s) := \int_{\mathbb{R}^N} G(t-s, y) u(s, y) dy.$$

Compute

$$g'(s) = \int -\Delta G u + G \Delta u + G u^{1+p} = \int G u^{1+p} \geq \left(\int G u \right)^{1+p} = g^{1+p}(s).$$

Based on this, find

$$f(t) \lesssim \frac{1}{t^{1/p}}.$$

Nonlocal dispersal (seeds)

$$\partial_t v = J * v - v.$$

J is a probability density on \mathbb{R}^N .

$J(x - y)$ is the probability of “jumping” from y to x .

$\int_{\mathbb{R}^N} J(x - y)v(t, y) dy = J * v(t, x)$ is the rate at which individuals arrive at x from all other positions.

$-\int_{\mathbb{R}^N} J(y - x)v(t, x) dy = -v(t, x)$ is the rate at which individuals leave x to reach any other positions.

Integro differential equations $\partial_t u = J * u - u + u^{1+p}$

Starting from a nontrivial compactly supported $u_0 \geq 0$, what happens to the Cauchy problem?

By nonlocal diffusion, individuals are sent in the region $u \approx 0$ where growth is not optimal (Allee effect) so it is more difficult to blow up.

The heavier the tails of J , the smaller should be the Fujita exponent...

On the linear equation $\partial_t v = J * v - v$

Assumption

$$\widehat{J}(\xi) = 1 - A|\xi|^\beta + o(|\xi|^\beta), \text{ as } \xi \rightarrow 0,$$

for some $0 < \beta \leq 2$, $A > 0$.

Rk: if J has a second momentum then $\beta = 2$, but **for heavier tails** $0 < \beta < 2$. In many cases β can be explicitly computed from the tails of J (see stable laws in probability theory).

Theorem (Chasseigne, Chaves and Rossi 2006)

Solutions to $\partial_t v = J * v - v$ decrease like $\frac{1}{(1+t)^{N/\beta}}$.

$$\partial_t u = J * u - u + u^{1+p}, \quad p > 0$$

Theorem (A. 2015)

Define $p_F := \frac{\beta}{N}$.

(i) $0 < p \leq p_F \implies$ all solutions blow up in finite time.

(ii) $p > p_F \implies$ some solutions (“small” u_0) are global and get extinct “like the Heat equation”.

Rk: as expected

$$p_F(\text{NONLOCAL}) \leq p_F(\text{LOCAL}).$$

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A general framework

We now work in the **half-space** \mathbb{R}_+^N and consider

$$\begin{cases} \partial_t u = Au + |u|^p u & t > 0, x \in \mathbb{R}_+^N, \\ u(0, x) = u_0(x) \geq 0 & x \in \mathbb{R}_+^N, \\ u(t, x) = 0 & t > 0, x \in \partial\mathbb{R}_+^N. \end{cases}$$

where $Au = \Delta u$ or $Au = -(-\Delta)^{\beta/2} u$ or $Au = J * u - u$.

► This Cauchy problem in \mathbb{R}_+^N is understood as the restriction of the Cauchy problem in \mathbb{R}^N obtained by anti-symmetrization.

Fujita critical exponent

Lemma (A., Kavian 2021)

Solutions to $\partial_t v = Av$ decrease like $\frac{1}{(1+t)^{(N+1)/\beta}}$.

Theorem (A., Kavian 2021)

$$p_F(\text{half-space}) = \frac{\beta}{N+1}.$$

Rk: as expected

$$p_F(\text{half-space}) \leq p_F(\text{whole space}).$$

Rk 1: where to evaluate v ?

The proof of the blow-up relies on a **lower bound** “with the good magnitude” of v the solution to the linear diffusion equation.

- ▶ In \mathbb{R}^N , considering $v(t, 0)$ was enough.
- ▶ In \mathbb{R}_+^N , we need to evaluate $v(t, \cdot)$ at an **appropriate moving point**:

Lemma (Pointwise estimate from below)

There exist two constants $\gamma > 0$ and $C = C(\gamma) > 0$ such that

$$v\left(t, \gamma t^{1/\beta} \mathbf{e}_N\right) \geq \frac{Cm_1(v_0)}{t^{(N+1)/\beta}}, \quad \forall t \gg 1.$$

Rk 2: role of the boundary conditions

Thus, in the half-space (say for $A = \Delta$):

$$p_F(\text{Dirichlet}) = \frac{2}{N+1}.$$

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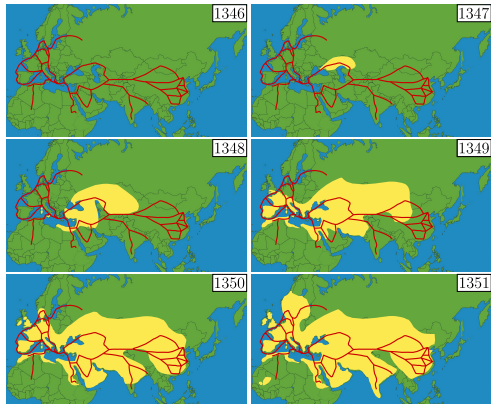
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Motivations

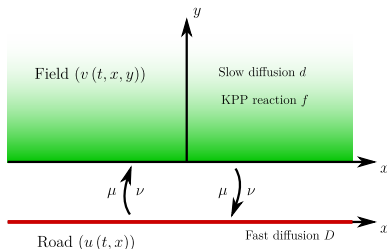


Motivations



The road-field model (Berestycki, Roquejoffre and Rossi)

$$\begin{cases} \partial_t v = d\Delta_{x,y}v + f_{KPP}(v), & t > 0, x \in \mathbb{R}^{N-1}, y > 0, \\ -d\partial_y v|_{y=0} = \mu u - \nu v|_{y=0}, & t > 0, x \in \mathbb{R}^{N-1}, \\ \partial_t u = D\Delta_x u + \nu v|_{y=0} - \mu u, & t > 0, x \in \mathbb{R}^{N-1}. \end{cases}$$



► $D > 2d \implies$ acceleration of invasion.

The purely diffusive road-field model

$$\begin{cases} \partial_t v = d\Delta_{x,y} v, & t > 0, x \in \mathbb{R}^{N-1}, y > 0, \\ -d\partial_y v|_{y=0} = \mu u - \nu v|_{y=0}, & t > 0, x \in \mathbb{R}^{N-1}, \\ \partial_t u = D\Delta_x u + \nu v|_{y=0} - \mu u, & t > 0, x \in \mathbb{R}^{N-1}. \end{cases}$$

Movie

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Movie

- ▶ Fourier (on the road variable)/Laplace (on time) transform is a good strategy and provides:

The solution explicitly

Theorem (A., Ducasse, Tréton 2022)

$$v(t, X) = \mathbb{V}(t, X) + \frac{\mu}{\sqrt{d}} \int_{\mathbb{R}^{N-1}} \Lambda(t, z, y) u_0(x - z) dz \\ + \frac{\mu \nu}{\sqrt{d}} \int_0^t \int_{\mathbb{R}^{N-1}} \Lambda(s, z, y) \mathbb{V}|_{y=0}(t - s, x - z) dz ds,$$

$$u(t, x) = e^{-\mu t} \mathbb{U}(t, x) \\ + \nu \int_0^t e^{-\mu(t-s)} \int_{\mathbb{R}^{N-1}} G_R(t - s, x - z) v|_{y=0}(s, z) dz ds,$$

where \mathbb{V} , \mathbb{U} , G_R are well-known while Λ , the keystone for writing the solution, is explicit but “not so nice”...

$\Lambda = \Lambda(t, x, y)$ is defined as

$$\frac{e^{-\frac{y^2}{4dt}}}{(2\pi)^{N-1}} \int_{\mathbb{R}^{N-1}} \left[a\alpha\Phi_\alpha + b\beta\Phi_\beta + c\gamma\Phi_\gamma \right] (t, \xi, y) e^{-dt\|\xi\|^2 + i\xi \cdot x} d\xi,$$

with $(\alpha, \beta, \gamma) = (\alpha, \beta, \gamma)(\xi)$ being the three complex roots of the δ -indexed polynomials

$$P_\delta(\sigma) = \sigma^3 + \frac{\nu}{\sqrt{d}}\sigma^2 + (\mu + \delta)\sigma + \frac{\nu\delta}{\sqrt{d}}, \quad \text{with } \delta = (D - d)\|\xi\|^2,$$

$(a, b, c) = (a, b, c)(\xi)$ being given by

$$a = \frac{1}{(\alpha - \beta)(\alpha - \gamma)}, \quad b = \frac{1}{(\beta - \alpha)(\beta - \gamma)}, \quad c = \frac{1}{(\gamma - \alpha)(\gamma - \beta)},$$

and for $\bullet \in \{\alpha, \beta, \gamma\}$,

$$\Phi_\bullet(t, \xi, y) = \frac{\text{Erfc}}{\Gamma} \left(\frac{-2\bullet\sqrt{dt} + y}{2\sqrt{dt}} \right),$$

where $\Gamma(\ell) = e^{-\ell^2}$, and Erfc is the complementary error function.

Asymptotic decay

Theorem (A., Ducasse, Tréton 2022)

We have

$$\|v(t, \cdot)\|_{L^\infty(\mathbb{R}_+^N)} \lesssim \frac{C_{v_0} \ln(1+t) + C_{u_0, v_0}}{(1+t)^{N/2}}, \quad \forall t > 0,$$

$$\|u(t, \cdot)\|_{L^\infty(\mathbb{R}^{N-1})} \lesssim \frac{C_{v_0} \ln(1+t) + C_{u_0, v_0}}{(1+t)^{N/2}}, \quad \forall t > 0.$$

Fujita blow-up phenomena on the road-field

$$\begin{cases} \partial_t v = d\Delta_{x,y}v + v^{1+p}, & t > 0, x \in \mathbb{R}^{N-1}, y > 0, \\ -d\partial_y v|_{y=0} = \mu u - \nu v|_{y=0}, & t > 0, x \in \mathbb{R}^{N-1}, \\ \partial_t u = D\Delta_x u + \nu v|_{y=0} - \mu u, & t > 0, x \in \mathbb{R}^{N-1}. \end{cases}$$

► Ongoing work by **Samuel Tréton**...

Thanks for your attention.