Numerical study of singular behavior in compressible flows

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Collaborators

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Main question:

Can vacuum appear in a compressible Navier-Stokes fluid?

or

Does there exist a weak solution to Navier-Stokes where the density ρ reaches zero in finite time, assuming $\rho(\cdot, 0)$ bounded away from zero?

Why should you care?

- open problem
- need to clarify notion of solution
- inviscid case is understood (Euler)
- cool numerical problem

Present work: numerical study of this theoretical question.

Disclaimer: No attempt is made at modeling interstellar matter and/or low density fluids.

Simplifications and notation

Symmetric flows, no swirl:

- $\rho(\mathbf{x}, t) = \rho(\mathbf{r}, t)$: density
- $\vec{u}(x,t) = \frac{x}{r} u(r,t)$: velocity

x point in space, r = |x|, t is time

This talk: barotropic flows: pressure depends solely on ρ

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Navier-Stokes

$$\rho_t + (\rho u)_{\xi} = 0 \qquad \text{mass}$$

$$\rho(u_t + uu_r) + \frac{1}{\gamma M^2} (\rho^{\gamma})_r = \frac{1}{\text{Re}} u_{\xi r} \qquad \text{momentum}$$

where

- M Mach number
- ► *Re* Reynolds number
- γ adiabatic coefficient
- $\blacktriangleright \ \partial_{\xi} = \partial_r + \frac{n-1}{r}$
- *n* spatial dimension (n = 1,2,3)

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Euler

Inviscid fluid: $\textit{Re} \rightarrow \infty$

$$\rho_t + (\rho u)_{\xi} = 0 \qquad \text{mass}$$

$$\rho(u_t + uu_r) + \frac{1}{\gamma M^2} (\rho^{\gamma})_r = 0 \qquad \text{momentum}$$

Riemann data (r > 0)

 $\begin{cases} \rho(r,0) &= 1, \\ u(r,0) &= 1, \end{cases} \quad 1D: \quad u(r,0) = \begin{cases} -1 & \text{if } r < 0, \\ 1 & \text{if } r > 0. \end{cases}$

"strength of the pull" is measured by M

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Riemman solution $M > \frac{2}{\gamma-1}$

$$\begin{bmatrix} \rho \\ \mu \end{bmatrix} (r,t) = \begin{cases} \begin{bmatrix} 1 \\ -1 \end{bmatrix} & \text{if } \frac{r}{t} < -1 - \frac{1}{M}, \\ \begin{bmatrix} \left(\frac{2}{\gamma+1} - M\frac{\gamma-1}{\gamma+1}(1+\frac{r}{t})\right)^{2/(\gamma-1)} \\ \frac{1}{M(\gamma+1)} \left(2 + (1-\gamma)M + 2M\frac{r}{t}\right) \end{bmatrix} & \text{if } -1 - \frac{1}{M} < \frac{r}{t} < -1 + \frac{2}{\gamma-1}\frac{1}{M}, \\ \begin{bmatrix} 0 \\ \emptyset \end{bmatrix} & \text{if } -1 + \frac{2}{\gamma-1}\frac{1}{M} < \frac{r}{t} < 1 - \frac{2}{\gamma-1}\frac{1}{M}, \\ \begin{bmatrix} \left(\frac{2}{\gamma+1} + M\frac{\gamma-1}{\gamma+1}(-1+\frac{r}{t})\right)^{2/(\gamma-1)} \\ \frac{1}{M(\gamma+1)} \left(-2 + (-1+\gamma)M + 2M\frac{r}{t}\right) \end{bmatrix} & \text{if } 1 - \frac{2}{\gamma-1}\frac{1}{M} < \frac{r}{t} < 1 + \frac{1}{M}, \\ \begin{bmatrix} 1 \\ 1 \end{bmatrix} & \text{if } 1 + \frac{1}{M} < \frac{r}{t}. \end{cases}$$

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Riemman solution $0 < M < \frac{2}{\gamma - 1}$

$$\begin{bmatrix} \rho \\ u \end{bmatrix} (r,t) = \begin{cases} \begin{bmatrix} 1 \\ -1 \end{bmatrix} & \text{if } \frac{r}{t} < -1 - \frac{1}{M}, \\ \begin{bmatrix} \left(\frac{2}{\gamma+1} - M\frac{\gamma-1}{\gamma+1}(1+\frac{r}{t})\right)^{2/(\gamma-1)} \\ \frac{1}{M(\gamma+1)}\left(2 + (1-\gamma)M + 2M\frac{r}{t}\right) \end{bmatrix} & \text{if } -1 - \frac{1}{M} < \frac{r}{t} < -\frac{1}{M} + \frac{\gamma-1}{2}, \\ \begin{bmatrix} \left(1 - \frac{M}{2}(\gamma-1)\right)^{\frac{2}{\gamma-1}} \\ 0 \end{bmatrix} & \text{if } -\frac{1}{M} + \frac{\gamma-1}{2} < \frac{r}{t} < \frac{1}{M} - \frac{\gamma-1}{2}, \\ \begin{bmatrix} \left(\frac{2}{\gamma+1} + M\frac{\gamma-1}{\gamma+1}(-1+\frac{r}{t})\right)^{2/(\gamma-1)} \\ \frac{1}{M(\gamma+1)}\left(-2 + (-1+\gamma)M + 2M\frac{r}{t}\right) \end{bmatrix} & \text{if } \frac{1}{M} - \frac{\gamma-1}{2} < r/t < 1 + \frac{1}{M}, \\ \begin{bmatrix} 1 \\ 1 \end{bmatrix} & \text{if } 1 + \frac{1}{M} < \frac{r}{t}. \end{cases}$$

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Known Euler results: 1D

Explicit Riemann solution: vacuum $\Leftrightarrow M > \frac{2}{\gamma-1}$



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2, 3 D "Riemann problem"

$$\rho = \rho(s), \qquad u = u(s), \qquad s = \frac{t}{r}$$

$$\Rightarrow \qquad \rho_s = (n-1)\frac{\rho u(1-su)}{s^2 c^2 - (1-su)^2}$$

$$u_s = (n-1)\frac{sc^2 u}{s^2 c^2 - (1-su)^2}$$
where $\rho(0) = 1, u(0) = 1, c = \frac{1}{M}\rho^{\frac{\gamma-1}{2}}$

Phase space analysis (Zheng, 2001) shows existence of critical Mach number M^*

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Known Euler results: 2, 3 D

Zheng (2001): vacuum $\Leftrightarrow M > M^*$



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Euler: phase diagram



(s = t/r)

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Euler: critical Mach number



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Known Navier-Stokes results

- ► theory far from complet, Danchin (2005), Feireisl (2004), Hoff (1997), P.L. Lions (1998)
- ► unique uniqueness result for discont. sol., Hoff (2006)
- Hoff & Smoller (2001): no vacuum formation for 1D NS
- Xin & Yuan (2006): 2, 3D sufficient conditions to rule out vacuum
- results below are consistent with the above

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Numerics

equations are split

$$(\rho_n, u_n) \xrightarrow{\text{Euler}} (\rho^*, u^*) \xrightarrow{\rho^* u_t = \frac{1}{R_\theta} u_{\xi r}} (\rho_{n+1}, u_{n+1})$$

- diffusive step solved by Chebyshev-Gauss-Radau collocation (avoid coord. singularity at 0)
- ► diffusive step advanced in time by BDF (can manage "infinite stiffness" when *ρ* = 0, i.e., index 1 DAE)
- Euler step advanced at each collocation node "à la Zheng" (ODE in s = t/r)

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Digression on collocation

Basic collocation principles

- Work on a finite grid
- ▶ Find *p* such that $p(x_j) = u_j$, $\forall x_j \in \text{grid}$
- approximate derivative is $p'(x_j)$.

Non periodic problems

- algebraic polynomials on non-uniform grids
- Chebyshev T_N optimality
- *T_N(x)* = cos(*N*θ) with θ = arccos x inherits fast convergence from periodic case

Coordinate singularity at r = 0

Chebyhsev-Gauss-Radau

(Spatial) discretization

• $u_N(r, t) = \sum_{i=0}^{N-1} U_i(t)\psi_i(r); \psi_i$ Lagrange interpolation pol.

The mesh



Euler vs NS, 3D, M = 1.2/2.7, Re = 10^6



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result: 3D



criterion: $\rho_N(r_{N-1}, t) < tol = 10^{-14}$, for some t, 0 < t < .005

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numerics \Rightarrow possible vacuum formation for multi-D NS flows

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Another example: relativistic Euler (2D)

$$\partial_t \hat{\rho} + \partial_x (\tilde{\rho} v_1) + \partial_y (\tilde{\rho} v_2) = 0,$$

$$\partial_t (\tilde{\rho} v_1) + \partial_x (\tilde{\rho} v_1^2 + \frac{1}{\gamma M^2} \rho^{\gamma}) + \partial_y (\tilde{\rho} v_1 v_2) = 0,$$

$$\partial_t (\tilde{\rho} v_2) + \partial_x (\tilde{\rho} v_1 v_2) + \partial_y (\tilde{\rho} v_2^2 + \frac{1}{\gamma M^2} \rho^{\gamma}) = 0,$$

where

$$\begin{split} \bullet \quad \tilde{\rho} &= \frac{\rho + \frac{1}{\gamma} \frac{\beta^2}{M^2} \rho^{\gamma}}{1 - \beta^2 |v|^2}, \qquad \hat{\rho} &= \tilde{\rho} - \frac{\beta^2}{\gamma M^2} \rho^{\gamma}, \\ \bullet \quad \beta &= \frac{\bar{v}}{c}, \\ \bullet \quad \beta \to \mathbf{0} \Rightarrow \text{classical Euler} \end{split}$$

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Singularity formation

- blow up of smooth compactly supported perturbations of constant states Pan & Smoller (2006)
- type of singularity is unknown
 - shock formation
 - violation of subluminal condition
 - mass concentration
- numerical difficulty: relationship between conserved and physical variables is non trivial

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Preliminary results



shock formation; more to follow...

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Conclusions

- analyzed two phenomena of singularity formation in compressible fluids
- discussed corresponding numerical challenges
- provided "numerical answers" to two open questions

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