THE CAHN-HILLIARD EQUATION WITH A LOGARITHMIC POTENTIAL AND DYNAMIC BOUNDARY CONDITIONS

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Cahn-Hilliard system :

$$\begin{array}{l} \frac{\partial u}{\partial t} = \kappa \Delta w, \ \kappa > 0\\ w = -\alpha \Delta u + f(u), \ \alpha > 0 \end{array}$$

Equivalently :

$$\frac{\partial u}{\partial t} + \alpha \kappa \Delta^2 u - \kappa \Delta f(u) = 0$$

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Describes the phase separation process in a binary alloy : spinodal decomposition, coarsening

u : order parameter

w: chemical potential

 κ : mobility

 α : related to the surface tension at the interface

f : derivative of a double-well potential *F* Typical choice :

$$F(s) = \frac{1}{4}(s^2 - 1)^2$$

f(s) = s^3 - s

Thermodynamically relevant potential :

$$\begin{split} F(s) &= -\theta_0 s^2 + \theta_1 ((1+s) \ln(1+s) \\ + (1-s) \ln(1-s)) \\ f(s) &= -2\theta_0 s + \theta_1 \ln \frac{1+s}{1-s} \\ s &\in (-1,1), \ 0 < \theta_1 < \theta_0 \end{split}$$

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Remark : κ should more generally depend on u and degenerate :

$$\frac{\partial u}{\partial t} = \operatorname{div}(\kappa(u)\nabla u)$$

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 $\kappa(s) = 1 - s^2$

Restricts the diffusion process to the interfacial region Is observed when the movements of atoms are confined to this region Usual boundary conditions :

$$\frac{\partial w}{\partial \nu} = 0 \text{ on } \Gamma$$
$$\frac{\partial u}{\partial \nu} = 0 \text{ on } \Gamma$$

 $\Omega \subset \mathbb{R}^N, N \leq 3$: domain occupied by the material $\Gamma = \partial \Omega$

 ν : unit outer normal vector

Equivalently :

$$\frac{\partial u}{\partial \nu} = \frac{\partial \Delta u}{\partial \nu} = 0 \text{ on } \Gamma$$

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Regular potentials :

• Well-posedness, regularity : C.M. Elliott-S. Zheng, B. Nicolaenko-B. Scheurer, D. Li-C. Zhong, ...

• Existence of finite-dimensional attractors : B. Nicolaenko-B. Scheurer-R. Temam, D. Li-C. Zhong, ...

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• Convergence of solutions to steady states : S. Zheng, P. Rybka-K.-H. Hoffmann

Logarithmic (singular) potentials :

Main difficulty : prove that *u* remains in (-1, 1)

Remark : Not true for regular potentials

• Well-posedness, regularity : C.M. Elliott-S. Luckhaus, C.M. Elliott-H. Garcke, A. Debussche-L. Dettori, A. Miranville-S. Zelik

• Existence of finite-dimensional attractors : A. Debussche-L. Dettori, A. Miranville-S. Zelik

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• Convergence of solutions to steady states : H. Abels-M. Wilke

Dynamic boundary conditions :

Influence of the walls for confined systems

Mainly studied for polymer mixtures

Technological applications

Problem : define the boundary conditions (we need 2 boundary conditions) First boundary condition : no mass flux at the boundary :

$$\frac{\partial w}{\partial \nu} = 0 \text{ on } \Gamma$$

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Different approach : G.R. Goldstein-A. Miranville-G. Schimperna

Second boundary condition : we consider, in addition to the Ginzburg-Landau free energy

$$\Psi_{\rm GL}(u,\nabla u) = \int_{\Omega} (\frac{\alpha}{2} |\nabla u|^2 + F(u)) dx$$

the surface free energy

$$\Psi_{\Gamma}(u,\nabla u) = \int_{\Gamma} (\frac{\alpha_{\Gamma}}{2} |\nabla_{\Gamma} u|^2 + G(u)) dx$$

 $\alpha_{\Gamma} > 0$ ∇_{Γ} : surface gradient

Original surface potential : $G(s) = \frac{1}{2}a_{\Gamma}s^2 - b_{\Gamma}s$

 $a_{\Gamma} > 0$: accounts for a modification of the effective interaction between the components

 b_{Γ} : characterizes the preferential attraction of one of the components by the walls

Total energy : $\Psi = \Psi_{GL} + \Psi_{\Gamma}$

The system tends to minimize the excess surface energy :

$$\frac{1}{d}\frac{\partial u}{\partial t} - \alpha_{\Gamma}\Delta_{\Gamma}u + g(u) + \alpha\frac{\partial u}{\partial\nu} = 0 \text{ on } \Gamma$$

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d > 0 : relaxation parameter Δ_{Γ} : Laplace-Beltrami operator g = G'

 \rightarrow Dynamic boundary condition

Regular potentials : the system is well understood

Contributors : R. Chill, C.G. Gal, E. Fašangová, A. Miranville, J. Pruess, R. Racke, H. Wu, S. Zelik, S. Zheng, ...

Singular potentials : more complicated and less understood

First existence and uniqueness result : G. Gilardi-A. Miranville-G. Schimperna

For f singular and g regular : sign assumptions on g near the singular points of f :

$$g(1) > 0, g(-1) < 0$$

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Forces the order parameter to stay away from ± 1 on Γ

Question :

• What happens when the sign conditions are not satisfied ?

Nonexistence of classical solutions :

When the sign conditions are not satisfied, we can have nonexistence of classical solutions

We consider the scalar ODE

$$y'' - f(y) = 0, x \in (-1, 1)$$

 $y'(\pm 1) = K > 0$

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Assumptions :

- $\bullet f$ is singular at ± 1
- $\bullet \; F(\pm 1) < +\infty \; (F'=f)$
- $\bullet f$ is odd

Satisfied by the usual logarithmic potentials

When *K* is small : existence and uniqueness of a solution which is separated from the singular values $(||y||_{L^{\infty}(-1,1)} < 1)$ and is odd

Standard interior regularity estimates yield

$$|y'(x)| \le c_0, \ |y(x)| \le 1 - \delta$$

 $x \in (-\frac{1}{2}, \frac{1}{2}), \delta > 0, c_0$ independent of *K*

Multiply the equation by y' and integrate over (0, 1):

$$|\frac{1}{2}K^2 - F(y(1))| \le c_1$$

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 c_1 (and $F(\pm 1)$) independent of K

This inequality cannot hold when K is large

 \rightarrow We do not have a classical solution

Since *y* is odd, the ODE can be rewritten as

$$y'' - f(y) = \langle y'' - f(y) \rangle$$

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 $<.>=\frac{1}{\operatorname{Vol}(\cdot)}\int_{\Omega}.dx$

 \rightarrow 1-D stationary Cahn-Hilliard system with dynamic BCs

Numerical results :

 $\Omega = (0,10)\times(0,4)$

Periodicity in the *x*-direction, dynamic boundary conditions in the *y*-direction u_0 : uniformly distributed random fluctuations of amplitude ± 0.5

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 $f(s) = -3s + \ln(\frac{1+s}{1-s}), g$ affine

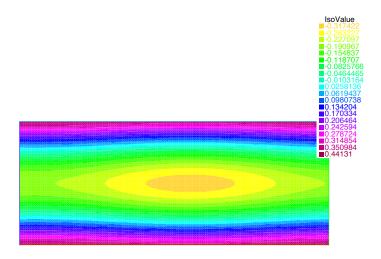
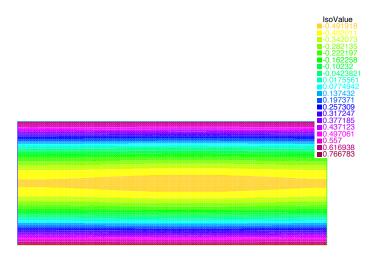


FIGURE: Isovalues of u, at time t = 20, g(s) = s - 0.8.

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FIGURE: Isovalues of u, at time t = 20, g(s) = s - 1.5.

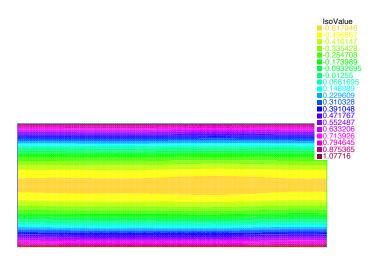


FIGURE: Isovalues of u, at time t = 0.72, g(s) = s - 3.

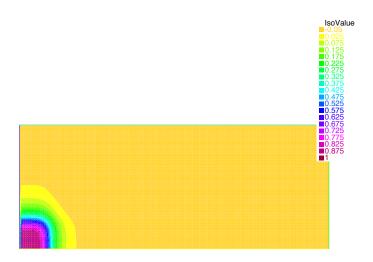


FIGURE: Isovalues of u_0 .

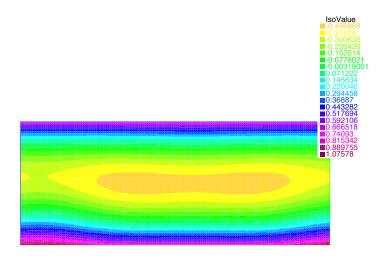


FIGURE: Isovalues of u, at time t = 0.46, g(s) = s - 3.

Convergence of a sequence of solutions to regularized problems :

$$\begin{aligned} \frac{\partial u}{\partial t} &= \Delta w \\ w &= -\Delta u + f_0(u) + \lambda u, \ \lambda \in R \\ \frac{\partial w}{\partial \nu} &= 0 \text{ on } \Gamma \\ \frac{\partial \psi}{\partial t} - \Delta_{\Gamma} \psi + g_0(\psi) + \psi + \frac{\partial u}{\partial \nu} = 0 \text{ on } \Gamma \\ \psi &= u|_{\Gamma} \end{aligned}$$

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$$f(s) = f_0(s) + \lambda s, \ g(s) = g_0(s) + s$$

Assumptions :

•
$$f_0 \in C^2(-1, 1), f_0(0) = 0$$

• $\lim_{s \to \pm 1} f_0(s) = \pm \infty, \lim_{s \to \pm 1} f'_0(s) = +\infty$
• $f'_0 \ge 0, \operatorname{sgn}(s) f''_0(s) \ge 0$
• $g_0 \in C^2(R), \|g_0\|_{C^2(R)} \le c$

Regularized potential :

$$\begin{aligned} f_{0,n}(s) &= f_0(s), \ |s| \le 1 - \frac{1}{n} \\ f_{0,n}(s) &= f_0(1 - \frac{1}{n}) + f_0'(1 - \frac{1}{n})(s - 1 + \frac{1}{n}) \\ s > 1 - \frac{1}{n} \\ f_{0,n}(s) &= f_0(-1 + \frac{1}{n}) + f_0'(-1 + \frac{1}{n})(s + 1 - \frac{1}{n}) \\ s < -1 + \frac{1}{n} \end{aligned}$$

Regularized problem : f_0 replaced by $f_{0,n}$

Existence and uniqueness of the solution u_n to the regularized problem

Satisfies, for *n* large enough

 $\Omega_{\epsilon} =$ $D_{\tau}u_n$

$$\begin{split} \|u_{n}(t)\|_{\mathcal{C}^{\alpha}(\Omega)}^{2} + \|u_{n}(t)\|_{H^{2}(\Gamma)}^{2} + \|u_{n}(t)\|_{H^{2}(\Omega_{\epsilon})}^{2} + \|u_{n}(t)\|_{H^{1}(\Omega)}^{2} + \\ \|\frac{\partial u_{n}}{\partial t}(t)\|_{H^{-1}(\Omega)}^{2} + \|\frac{\partial u_{n}}{\partial t}(t)\|_{L^{2}(\Gamma)}^{2} + \\ \|\nabla D_{\tau}u_{n}(t)\|_{L^{2}(\Omega)}^{2N} + \|f_{0,n}(u_{n}(t))\|_{L^{1}(\Omega)} + \\ \int_{t}^{t+1} (\|\frac{\partial u_{n}}{\partial t}(s)\|_{H^{-1}(\Omega)}^{2} + \|\frac{\partial u_{n}}{\partial t}(s)\|_{L^{2}(\Gamma)}^{2}) ds \leq \\ ce^{-\beta t}(1 + \|u_{n}(0)\|_{H^{1}(\Omega)}^{2} + \|u_{n}(0)\|_{H^{1}(\Gamma)}^{2} + \\ \|\frac{\partial u_{n}}{\partial t}(0)\|_{H^{-1}(\Omega)}^{2} + \|\frac{\partial u_{n}}{\partial t}(0)\|_{L^{2}(\Gamma)}^{2})^{2} + c' \\ \Omega_{\epsilon} = \{x \in \Omega, \ d(x, \Gamma) > \epsilon\}, \ \epsilon > 0 \\ D_{\tau}u_{n} = \nabla u_{n} - \frac{\partial u_{n}}{\partial \nu}\nu \\ \alpha > 0, \ \beta > 0, \ c, \ c' \ \text{independent of } n \end{split}$$

Remark : Actually, $u_n(t) \in H^2(\Omega)$, but this regularity does not pass to the limit

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Smoothing property :

$$\begin{aligned} \|\frac{\partial u_n}{\partial t}(t)\|_{H^{-1}(\Omega)}^2 + \|\frac{\partial u_n}{\partial t}(t)\|_{L^2(\Gamma)}^2 \leq \\ \frac{c}{t}(1+\|u_n(0)-\langle u_n(0)\rangle \|_{H^{-1}(\Omega)}^2 + \|u_n(0)\|_{L^2(\Gamma)}^2) \end{aligned}$$

 $t \in (0, 1], c$ independent of n

Lipschitz estimate :

$$\begin{split} \|u_1(t) - u_2(t)\|_{H^{-1}(\Omega)} + \\ \|u_1(t) - u_2(t)\|_{L^2(\Gamma)} &\leq \\ ce^{c't}(\|u_1(0) - u_2(0)\|_{H^{-1}(\Omega)} + \\ \|u_1(0) - u_2(0)\|_{L^2(\Gamma)}) \\ &< u_1(0) > = < u_2(0) > = m, \ t \ge 0 \end{split}$$

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c, c' independent of t, n, u_1, u_2

 u_n converges to some function u

We wish to call *u* the "generalized" solution to the singular problem Variational solutions :

We set

$$\begin{split} B(u,v) &= (\nabla u, \nabla v)_{\Omega} + \lambda(u,v)_{\Omega} + \\ + L((-\Delta)^{-1}\overline{u}, \overline{v})_{\Omega} + (\nabla_{\Gamma} u, \nabla_{\Gamma} v)_{\Gamma} \end{split}$$

 $u, v \in H^1(\Omega) \otimes H^1(\Gamma) = \{w, w \in H^1(\Omega), w|_{\Gamma} \in H^1(\Gamma)\}$

L > 0 chosen s.t.

$$\begin{split} \|\nabla u\|_{L^2(\Omega)^3}^2 + \lambda \|u\|_{L^2(\Omega)}^2 + L\|u\|_{H^{-1}(\Omega)}^2 \geq \\ \frac{1}{2} \|u\|_{H^1(\Omega)}^2, \ u \in H^1(\Omega), \ < u >= 0 \end{split}$$

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 $\overline{u} = u - \langle u \rangle$ (.,.)_{\Omega}, (.,.)_{\Gamma} : scalar products in $L^2(\Omega)$ and $L^2(\Gamma)$ We rewrite the problem as

$$(-\Delta)^{-1}\frac{\partial u}{\partial t} - \Delta u + f_0(u) + \lambda u - \langle w \rangle = 0$$

$$w = -\Delta u + f_0(u) + \lambda u$$

$$\frac{\partial \psi}{\partial t} - \Delta_{\Gamma} \psi + g(\psi) + \frac{\partial u}{\partial \nu} = 0 \text{ on } \Gamma$$

$$\psi = u|_{\Gamma}$$

$$u|_{t=0} = u_0, \ \psi|_{t=0} = \psi_0$$

We multiply the first equation by u - v, v = v(x) s.t.

$$< u(t) - v >= 0, t \ge 0:$$

$$((-\Delta)^{-1}\frac{\partial u}{\partial t}, u - v)_{\Omega} + (\frac{\partial u}{\partial t}, u - v)_{\Gamma} + B(u, u - v) + (f_0(u), u - v)_{\Omega} = L(u, (-\Delta)^{-1}(u - v))_{\Omega} - (g(u), u - v)_{\Gamma}$$

Positivity of B and monotonicity of f_0 :

$$((-\Delta)^{-1}\frac{\partial u}{\partial t}, u-v)_{\Omega} + (\frac{\partial u}{\partial t}, u-v)_{\Gamma} + B(v, u-v) + (f_0(v), u-v)_{\Omega} \leq L(u, (-\Delta)^{-1}(u-v))_{\Omega} - (g(u), u-v)_{\Gamma}$$

Variational inequality (VI)

We set

$$\Phi = \{ (u, \psi) \in L^{\infty}(\Omega) \times L^{\infty}(\Gamma), \\ \|u\|_{L^{\infty}(\Omega)} \le 1, \|\psi\|_{L^{\infty}(\Gamma)} \le 1 \}$$

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Definition : Let $(u_0, \psi_0) \in \Phi$. Then, (u, ψ) is a variational solution if

(i)
$$u(t)|_{\Gamma} = \psi(t)$$
 a.e. $t > 0$, $u(0) = u_0$, $\psi(0) = \psi_0$;
(ii) $-1 < u(t, x) < 1$ a.e. $(t, x) \in \mathbb{R}^+ \times \Omega$;
(iii) $(u, \psi) \in \mathcal{C}([0, +\infty); H^{-1}(\Omega) \times L^2(\Gamma)) \cap L^2(0, T; H^1(\Omega) \times H^1(\Gamma)),$
 $T > 0$;

(iv)
$$f(u) \in L^1((0,T) \times \Omega), T > 0;$$

(v) $\left(\frac{\partial u}{\partial t}, \frac{\partial \psi}{\partial t}\right) \in L^2(\tau, T; H^{-1}(\Omega) \times L^2(\Gamma)), T > \tau > 0;$
(vi) $\langle u(t) \rangle = \langle u_0 \rangle, t \ge 0;$

(vii) the variational inequality (VI) is satisfied for a.e. t > 0 and every test function v = v(x) s.t. $v \in H^1(\Omega) \otimes H^1(\Gamma)$, $f(v) \in L^1(\Omega)$, $\langle v \rangle = \langle u_0 \rangle$.

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Remark : $u(t)|_{\Gamma} = \psi(t)$ only for t > 0

• A variational solution, if it exists is unique

• $\forall (u_0, \psi_0) \in \Phi, \exists$ a variational solution and $(u_n, \psi_n = u_n|_{\Gamma})$ converges (for a subsequence) to a variational solution

- The variational solutions satisfy the a priori estimates mentioned earlier
- The variational solutions satisfy the smoothing and Lipschitz properties

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A variational solution does not necessarily solve the equations in the usual sense

True if $u(t) \in H^2(\Omega)$

A variational solution solves the first equation

$$(-\Delta)^{-1} \frac{\partial u}{\partial t} - \Delta u + f_0(u) + \lambda u - \langle w \rangle = 0 \text{ in } \mathcal{D}'$$

Does not necessarily satisfy the dynamic boundary condition

More precisely, the trace

$$\frac{\partial u}{\partial \nu} = \left[\frac{\partial u}{\partial \nu}\right]_{\text{int}}$$

exists in $L^{\infty}(\tau, T; L^1(\Gamma)), 0 < \tau < T$

 (u_n, ψ_n) satisfies

$$\frac{\partial \psi_n}{\partial t} - \Delta_{\Gamma} \psi_n + g(\psi_n) + \frac{\partial u_n}{\partial \nu} = 0 \text{ on } \Gamma$$

in $L^{\infty}(\tau, T; L^{2}(\Gamma)), T > \tau > 0$, and the limit

$$\left[\frac{\partial u}{\partial \nu}\right]_{\text{ext}} = \lim_{n \to +\infty} \frac{\partial u_n}{\partial \nu}$$

exists in $L^{\infty}(\tau, T; L^2(\Gamma))$ weak star

$$\rightarrow \frac{\partial \psi}{\partial t} - \Delta_{\Gamma} \psi + g(\psi) + \left[\frac{\partial u}{\partial \nu}\right]_{\text{ext}} = 0 \text{ on } \Gamma$$

 \rightarrow A variational solution is a classical one if

$$\left[\frac{\partial u}{\partial \nu}\right]_{\text{int}} = \left[\frac{\partial u}{\partial \nu}\right]_{\text{ext}} \text{ a.e. } (t, x) \in \mathbb{R}^+ \times \Gamma$$

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Remark : Scalar ODE

$$y'' - f(y) = 0, x \in (-1, 1)$$

 $y'(\pm 1) = K > 0$

There exists a critical value K_0 s.t., if $K > K_0$, there is no classical solution

However, there exists a variational solution which is solution to

$$y'' - f(y) = 0, x \in (-1, 1)$$

 $y(\pm 1) = \pm 1$
 $y'|_{x=\pm 1} \neq K$

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Existence of classical solutions :

Related to the H^2 -regularity and the separation from the singularities of f_0 **Theorem :** Let (u, ψ) be a variational solution and set, for $\delta > 0$ and T > 0,

$$\Omega_{\delta}(T) = \{ x \in \Omega, \ |u(T,x)| < 1 - \delta \}.$$

Then, $u(T) \in H^2(\Omega_{\delta}(T))$ and

$$\|u(T)\|_{H^2(\Omega_{\delta}(T))} \leq Q_{\delta,T},$$

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where $Q_{\delta,T}$ is independent of u.

Consequence : if

$$|u(t,x)| < 1$$
 a.e. $(t,x) \in \mathbb{R}^+ \times \Gamma$

then

$$\left[\frac{\partial u}{\partial \nu}\right]_{\text{int}} = \left[\frac{\partial u}{\partial \nu}\right]_{\text{ext}} \text{ a.e. } (t, x) \in \mathbb{R}^+ \times \Gamma$$

and u is a classical solution

 \rightarrow The existence of classical solutions is related to the separation property on the boundary

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True if f_0 has sufficiently strong singularities

Theorem : We assume that

$$\lim_{s \to \pm 1} F_0(s) = +\infty, \ F'_0 = f_0.$$

Then, the separation property on the boundary holds and a variational solution is a classical one.

True if f_0 behaves like $\frac{s}{(1-s^2)^p}$, p > 1

Not true for logarithmic potentials

In that case, we can have |u(t,x)| = 1 on a set with nonzero measure on the boundary (possibly, on the whole boundary)

Theorem : We assume that

$$\pm g(\pm 1) > 0.$$

Then, a variational solution is a classical one.

Existence of finite-dimensional attractors :

Conservation of the total mass $(\langle u \rangle)$: we restrict ourselves to

$$\Phi_m = \{(u, \psi) \in \Phi, \ < u >= m\}, \ m \in (-1, 1)$$

Theorem : For every $m \in (-1, 1)$, the semigroup S(t) acting on Φ_m possesses the finite-dimensional global attractor \mathcal{A}_m (in $H^{-1}(\Omega) \times L^2(\Gamma)$) which is bounded in $\mathcal{C}^{\alpha}(\Omega) \times \mathcal{C}^{\alpha}(\Gamma)$, $0 < \alpha < \frac{1}{4}$.

Global attractor : unique compact set of Φ_m which is invariant $(S(t)A_m = A_m, t \ge 0)$ and attracts all bounded sets of initial data

Existence of the global attractor : follows from classical results

Finite-dimensionality : construction of an exponential attractor

Exponential attractor : compact and positively invariant $(S(t)\mathcal{M}_m \subset \mathcal{M}_m, t \ge 0)$ set which contains the global attractor and has finite fractal dimension

We need some (asymptotically) compact smoothing property on the difference of 2 solutions

We have

$$\begin{aligned} \|u_1(t) - u_2(t)\|_{\Phi^{w}}^2 &\leq c e^{-\beta t} \|u_1(0) - u_2(0)\|_{\Phi^{w}}^2 + c' \int_0^t \|\theta(u_1(s) - u_2(s))\|_{L^2(\Omega)}^2 ds \end{aligned}$$

 $\beta > 0, \theta$: smooth cut-off function

 $\Phi^{\mathrm{w}} = H^{-1}(\Omega) \times L^2(\Gamma)$

 \rightarrow The semigroup is a contraction, up to $\|\theta(u_1 - u_2)\|_{L^2(0,t;L^2(\Omega))}$

Compactness : We work on spaces of trajectories and use the compactness of

$$L^2(0,t;H^1(\Omega)) \cap H^1(0,t;H^{-3}(\Omega)) \subset L^2(0,t;L^2(\Omega))$$

We have

$$\begin{aligned} &\|\frac{\partial}{\partial t}[\theta(u_1 - u_2)]\|^2_{L^2(0,t;H^{-3}(\Omega))} + \\ &\|\theta(u_1 - u_2)\|^2_{L^2(0,t;H^{-1}(\Omega))} \leq \\ &ce^{c't}\|u_1(0) - u_2(0)\|^2_{H^{-1}(\Omega)\cap L^2(\Gamma)} \end{aligned}$$

 $u_1(0), u_2(0) \in B_{H^{-1}(\Omega) \cap L^2(\Gamma)}(u_0, \epsilon), \epsilon > 0$ small