A NONHOMOGENEOUS BOUNDARY VALUE PROBLEM IN MASS TRANSFER THEORY

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ABSTRACT. We prove a uniqueness result of solutions for a system of PDEs of Monge-Kantorovich type arising in problems of mass transfer theory. The results are obtained under very mild regularity assumptions both on the reference set $\Omega \subset \mathbb{R}^n$, and on the (possibly asymmetric) norm defined in Ω . In the special case when Ω is endowed with the Euclidean metric, our results provide a complete description of the stationary solutions to the tray table problem in granular matter theory.

1. Introduction

The model system usually considered for the description of the stationary configurations of sandpiles on a tray table is the Monge-Kantorovich type system of PDEs

$$\begin{cases}
-\operatorname{div}(v D u) = f & \text{in } \Omega, \\
|D u| \le 1, \ v \ge 0 & \text{in } \Omega, \\
(1 - |D u|)v = 0 & \text{in } \Omega, \\
u = \phi & \text{on } \partial\Omega
\end{cases}$$

(see, e.g., [4, 10, 23]). The data of the problem are the flat surface of the table $\Omega \subseteq \mathbb{R}^2$, the profile of the tray table ϕ , and the density of the source $f \geq 0$. The dynamical behavior of the granular matter is pictured by the pair (u,v), where u is the profile of the standing layer, whose slope has not to exceed a critical value $(|Du| \leq 1)$ in order to prevent avalanches, while $v \geq 0$ is the thickness of the rolling layer (see Figure 1). The condition (1-|Du|)v=0 corresponds to require that the matter runs down only in the region where the slope of the heaps is maximal. On the border of the table the standing layer fills the gap with the height ϕ , and the exceeding sand falls down. If $\phi=0$, the model reduces to the open table problem, largely investigated in recent years (see e.g. [10, 11, 12, 28] and the references therein).

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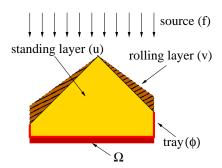


FIGURE 1. The tray table problem

In the present paper we consider a more general system of PDEs in the open, bounded and connected set $\Omega \subseteq \mathbb{R}^n$

(1)
$$\begin{cases} -\operatorname{div}(v D\rho(Du)) = f & \text{in } \Omega, \\ \rho(Du) \le 1, \ v \ge 0 & \text{in } \Omega, \\ (1 - \rho(Du))v = 0 & \text{in } \Omega, \\ u = \phi & \text{on } \partial\Omega, \end{cases}$$

in the unknowns $v \in L^1_+(\Omega)$, $u \in W^{1,\infty}(\Omega) \cap C(\overline{\Omega})$. The data are:

- $K \subseteq \mathbb{R}^n$, a compact convex set of class C^1 containing the origin in its interior, and its (strictly convex) polar set K^0 . The set K appears in (1) by means of its gauge function $\rho \colon \mathbb{R}^n \to [0, +\infty)$. Moreover the closure of Ω will be equipped with a (possibly asymmetric) geodesic distance d_L given in terms of the gauge function ρ^0 of the polar set K^0 (see Section 3);
- $f \in L^1_+(\Omega)$;
- $\phi: \partial\Omega \to \mathbb{R}^n$, 1-Lipschitz function w.r.t. d_L .

This general version takes into account the possibility of homogeneous anisotropies, and can be applied, for example, to the Bean's model for the description of the macroscopic electrodynamics of hard superconductors (see [16, 17, 18, 19], for the case $\phi = 0$ and Ω , K smooth sets), and to optimal shape design (see [8]).

S. Bianchini in [6] proved that the Lax-Hopf function u_{ϕ} (defined in (13) below), coupled with an explicit $v_f \in L^1_+(\Omega)$ solves (1) (see Theorem 6.9). The function v_f is constructed as follows. The set Ω can be covered by a family of disjoint transport rays, and every v-component of a solution to (1) satisfies a first order linear ODE along almost every ray. The function v_f is then uniquely defined along each ray, by requiring that it vanishes at the final point. This procedure gives a solution to (1), thanks to a disintegration formula for the Lebesgue measure along the rays.

In this paper we shall provide a complete characterization of the solutions (u, v) to (1). Starting from the fact that (u_{ϕ}, v_f) is a solution, we show that a pair (u, v) is a solution if and only if both (u, v_f) and (u_{ϕ}, v) are solutions, so that we can decouple the original problem (see Section 5). Moreover, the

admissible u-components are characterized as the solutions of a minimization problem which has an explicit minimal solution u_f depending on the source term f. Since u_{ϕ} is the maximal solution of the problem, the uniqueness of the u-component corresponds to the case when $u_f = u_{\phi}$. We prove that this can occur if and only if the final points in Ω of the transport rays are contained in the support of the source f (see Section 7). Concerning the v-component, in [6] it is proved that every admissible v has to vanish at the final points in Ω of the transport rays. Hence, by the very definition of v_f , it is the unique v-component if the union T of the rays with both endpoints on $\partial\Omega$ has null Lebesgue measure. We shall prove that the converse is also true: if the set T has positive Lebesgue measure, we can construct a nontrivial function $v \in L^1_+(\Omega)$ such that $-\operatorname{div}(v D\rho(Du_{\phi})) = 0$, so that the uniqueness of the v-component fails.

These results generalize those in [16, 17, 18] to the case when the reference set Ω is possibly non-regular, the boundary datum ϕ is possibly non-homogeneous, and the constraint K is a possibly non-strictly convex set. In this case we have to use a completely different strategy with respect to the previous approaches.

Namely, if Ω and K are smooth enough, then it is possible to define a suitable notion of (anisotropic) curvatures of $\partial\Omega$, and the analysis of solutions to (1) can be performed using the methods of Differential Geometry (see, e.g., [13, 16, 17, 18, 24, 25, 27]). In the non-smooth case these methods do not work, and a different approach is needed, based on the techniques developed in Optimal Transport Theory (see [1, 6, 31]). Namely, system (1) can be understood as the PDE formulation of an optimal transport problem with strictly convex cost function ρ^0 , assigned initial distribution $\mu = f dx$ and transport potential given by the Lax-Hopf function u_{ϕ} (so that the final mass distribution ν is the measure concentrated on $\partial\Omega$ induced by the transportation of μ along the rays associated to u_{ϕ}). From this point of view, the minimization problem considered in Section 5 turns out to be the dual formulation of the Monge problem (see, e.g., [31] for an exhaustive presentation of this subject).

A comment on the regularity assumption on K is in order. Namely, if K is of class C^1 , then its polar set K^0 need not be of class C^1 , but it is strictly convex. It is well-known that the strict convexity of K^0 , which is the unit ball of the dual norm ρ^0 , allows non-trivial simplifications in the proofs; for example, the direction of the transport rays coincides with $D\rho(Du)$ almost everywhere in Ω . Related results in the case of a general possibly non-regular constraint K can be found in [7].

The paper is organized as follows. In Section 2 we collect all the notation and preliminary results. In Section 3 we recall some generalization of the McShane's extension theorem for Lipschitz functions with constrained gradient. These general results will be used for introducing the maximal and the minimal extension of the boundary datum ϕ consistent with the gradient constraint. Section 4 is devoted to the statement of the problem, some remarks on the hypotheses, and the statement of the existence result

due to S. Bianchini (see [6, Thm. 7.1]). One of the main tools needed for proving the uniqueness results is the fact that (1) is the Euler-Lagrange necessary condition for a minimization problem with a gradient constraint. This property is proved in Section 5. The necessary and sufficient conditions for the uniqueness of the solutions (u, v) to (1) are proved in Sections 6 and 7. As a byproduct of our analysis, with a little additional effort, we provide a complete description of the singular sets considered in this kind of optimal transport problems.

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2. NOTATION AND PRELIMINARIES

The standard scalar product of $x, y \in \mathbb{R}^n$ will be denoted by $\langle x, y \rangle$, while |x| will denote the Euclidean norm of x. Concerning the segment joining x with y, we set

$$[x, y] := \{tx + (1 - t)y; t \in [0, 1]\}, \quad [x, y] := [x, y] \setminus \{x, y\}.$$

Given a set $A \subseteq \mathbb{R}^n$, its interior, its closure and its boundary will be denoted by int A, \overline{A} and ∂A respectively.

A bounded open set $\mathcal{O} \subset \mathbb{R}^n$ (or, equivalently, $\overline{\mathcal{O}}$ or $\partial \mathcal{O}$) is of class C^k , $k \in \mathbb{N}$, if for every point $x_0 \in \partial \mathcal{O}$ there exists a ball $B = B_r(x_0)$ and a one-to-one mapping $\psi \colon B \to D$ such that $\psi \in C^k(B)$, $\psi^{-1} \in C^k(D)$, $\psi(B \cap \mathcal{O}) \subseteq \{x \in \mathbb{R}^n; x_n > 0\}$, $\psi(B \cap \partial \mathcal{O}) \subseteq \{x \in \mathbb{R}^n; x_n = 0\}$.

Let us now fix the notation and the basic results concerning the convex set which plays the role of gradient constraint for the u-component in (1). In the following we shall assume that

(2) K is a compact, convex subset of \mathbb{R}^n of class C^1 , with $0 \in \text{int } K$.

Let us denote by K^0 the polar set of K, that is

$$K^0 := \{ p \in \mathbb{R}^n; \ \langle p, x \rangle \le 1 \ \forall x \in K \}.$$

We recall that, if K satisfies (2), then K^0 is a compact, strictly convex subset of \mathbb{R}^n containing the origin in its interior, and $K^{00} = (K^0)^0 = K$ (see, e.g., [30]).

The gauge function $\rho \colon \mathbb{R}^n \to \mathbb{R}$ of K is defined by

$$\rho(\xi) := \inf\{t \ge 0; \ \xi \in tK\} = \max\{\langle \xi, \eta \rangle, \ \eta \in K^0\}, \quad \xi \in \mathbb{R}^n.$$

It is straightforward to see that ρ is a positively 1-homogeneous convex function such that $K = \{\xi \in \mathbb{R}^n \colon \rho(\xi) \leq 1\}$. The gauge function of the set K^0 will be denoted by ρ^0 .

The properties of the gauge functions needed in the paper are collected in the following theorem.

Theorem 2.1. Assume that $K \subseteq \mathbb{R}^n$ satisfies (2). Then the following hold: (i) ρ is continuously differentiable in $\mathbb{R}^n \setminus \{0\}$, and

$$\rho(\xi + \eta) \le \rho(\xi) + \rho(\eta) \quad \forall \ \xi, \eta \in \mathbb{R}^n.$$

(ii) K^0 is strictly convex, and

$$\rho^{0}(\xi + \eta) \leq \rho^{0}(\xi) + \rho^{0}(\eta), \quad \forall \ \xi, \eta \in \mathbb{R}^{n}$$
$$\rho^{0}(\xi + \eta) = \rho^{0}(\xi) + \rho^{0}(\eta) \iff \exists \lambda \geq 0 : \ \xi = \lambda \eta \ or \ \eta = \lambda \xi.$$

(iii) For every $\xi \neq 0$, $D\rho(\xi)$ belongs to ∂K^0 , and

$$\langle D\rho(\xi), \xi \rangle = \rho(\xi), \quad \langle p, \xi \rangle < \rho(\xi) \ \forall p \in K^0, \ p \neq D\rho(\xi).$$

Proof. See [30], Section 1.7.

In what follows we shall consider \mathbb{R}^n endowed with the possibly asymmetric norm $\rho^0(x-y)$, $x,y\in\mathbb{R}^n$. By Theorem 2.1(ii), the unit ball K^0 of ρ^0 is strictly convex but, under the sole assumption (2), it need not be differentiable. Moreover, the Minkowski structure (\mathbb{R}^n, ρ^0) is not a metric space in the usual sense, since ρ^0 need not be symmetric (for an introduction to nonsymmetric metrics see [22]). Finally, since K^0 is compact and $0 \in \text{int } K^0$, then the convex metric is equivalent to the Euclidean one, that is there exist $c_1, c_2 > 0$ such that $c_1|\xi| \leq \rho^0(\xi) \leq c_2|\xi|$ for every $\xi \in \mathbb{R}^n$.

The main motivation for introducing the convex metric associated to ρ^0 is the fact that the Sobolev functions with the gradient constrained to belong to K are the locally 1-Lipschitz functions with respect to ρ^0 , as stated in the following result.

Lemma 2.2. Let $\mathcal{O} \subset \mathbb{R}^n$ be a nonempty open set, and assume that the set $K \subset \mathbb{R}^n$ satisfies (2). Let ρ^0 be the gauge function of K^0 . Then the following properties are equivalent.

- i) $u: \mathcal{O} \to \mathbb{R}$ is a locally 1-Lipschitz function with respect to ρ^0 , i.e.
- (3) $u(x_2) u(x_1) \le \rho^0(x_2 x_1)$ for every $[x_1, x_2] \subset \mathcal{O}$.
 - ii) $u \in W^{1,\infty}(\mathcal{O})$, and $Du(x) \in K$ for a.e. $x \in \mathcal{O}$.

Proof. Let u be a locally 1-Lipschitz function in \mathcal{O} with respect to the metric ρ^0 . Then u is locally Lipschitz in \mathcal{O} , and if x_0 is a differentiability point of u, by (3) we have

$$o(|t|) = \frac{u(x_0 + t\xi) - u(x_0) - t\langle Du(x_0), \xi \rangle}{t} \le \rho^0(\xi) - \langle Du(x_0), \xi \rangle$$

$$\le 1 - \langle Du(x_0), \xi \rangle, \qquad \forall \xi \in K^0.$$

Hence $\langle Du(x_0), \xi \rangle \leq 1$ for all $\xi \in K^0$, that is $Du(x_0) \in (K^0)^0 = K$. Since u is a locally Lipschitz function in \mathcal{O} with gradient in the bounded set K, (ii) holds.

Conversely, let u satisfy (ii), and let $[x_1, x_2] \subset \mathcal{O}$. Clearly, we can assume $x_1 \neq x_2$. Let $S_r := \{z \in B_r(0) : \langle z - x_1, x_2 - x_1 \rangle = 0\}$, and let $Q_r := \{\bigcup [x_1 + z, x_2 + z] : z \in S_r\}$ be the cylinder of radius r around the segment $[x_1, x_2]$. Let us choose r > 0 such that $Q_r \subset \Omega$. By (ii) there exists

a set $N \subset S_r$ with (n-1)-Lebesgue measure zero, such that for every $z \in S_r \setminus N$, u is absolutely continuous along the segment $[x_1 + z, x_2 + z]$, and $Du(x_1 + z + s(x_2 - x_1)) \in K$ for a.e. $s \in [0,1]$. Let $(z_j) \subset S_r \setminus N$, $z_j \to 0$. Then, for every $j \in \mathbb{N}$,

$$u(x_2+z_j)-u(x_1+z_j) = \int_0^1 \langle Du(x_1+z_j+s(x_2-x_1)), x_2-x_1 \rangle \, ds \le \rho^0(x_2-x_1).$$

The conclusion now follows from the continuity of u in \mathcal{O} , passing to the limit in j.

In the sequel we need to consider locally 1-Lipschitz functions with respect to ρ^0 that agree with a given continuous boundary datum. This can be done only if the boundary values are compatible with the gradient bounds, and, in absence of assumptions on the regularity of the reference set, the compatibility condition concerns the variations of the data along every admissible path.

Here we introduce some notation for the curves, while the compatibility condition will be discussed in Section 3.

Definition 2.3. Let $D \subset \mathbb{R}^n$ be a nonempty compact set. For every $x_1, x_2 \in D$, $\Gamma^D_{x_1,x_2}$ will denote the (possibly empty) family of absolutely continuous curves $\gamma \colon [0,1] \to D$ such that $\gamma(0) = x_1$ and $\gamma(1) = x_2$.

For every absolutely continuous curve $\gamma \colon [0,1] \to \mathbb{R}^n$, let us denote by $L(\gamma)$ its length with respect to the convex metric associated to ρ^0 , that is

$$L(\gamma) := \int_0^1 \rho^0(\gamma'(t)) dt.$$

If D is a compact subset of \mathbb{R}^n , by a standard compactness argument we have that if $\Gamma^D_{x_1,x_2} \neq \emptyset$, then there exists a geodesic in D joining x_1 to x_2 , i.e. there exists a curve $\tilde{\gamma} \in \Gamma^D_{x_1,x_2}$ such that $L(\tilde{\gamma}) \leq L(\gamma)$ for every $\gamma \in \Gamma^D_{x_1,x_2}$ (see e.g. [2, Thm. 4.3.2], [15, §14.1]).

For sake of completeness we briefly recall that any geodesic curve lying in the interior of D is in fact a segment, due to strict convexity of ρ^0 .

Lemma 2.4. Let $D \subset \mathbb{R}^n$ be a nonempty compact set. Let $x_1, x_2 \in D$ be two points such that $\Gamma^D_{x_1, x_2} \neq \emptyset$, and let $\gamma \in \Gamma^D_{x_1, x_2}$ be a geodesic. If $\gamma(t) \in \text{int } D$ for every $t \in (0, 1)$, then the support of γ is the segment $[x_1, x_2]$.

Proof. By an approximation argument we can assume that $x_1, x_2 \in \text{int } D$ so that $\gamma(t) \in \text{int } D$ for every $t \in [0,1]$. Let $\eta(t) = x_1 + t(x_2 - x_1)$, $t \in [0,1]$, be a parametrization of the segment $[x_1, x_2]$. Assume by contradiction that the support of γ is not the segment $[x_1, x_2]$. Since K^0 is strictly convex we have that $L(\gamma) > L(\eta) = \rho^0(x_2 - x_1)$. Since the support of γ is contained in int D, there exists $\varepsilon \in (0,1)$ such that the support of $\gamma_{\varepsilon} := \gamma + \varepsilon(\eta - \gamma)$ is contained in int D. Then

$$L(\gamma_{\varepsilon}) = \int_{0}^{1} \rho^{0}((1-\varepsilon)\gamma' + \varepsilon\eta') \leq \int_{0}^{1} [(1-\varepsilon)\rho^{0}(\gamma') + \varepsilon\rho^{0}(\eta')]$$

= $(1-\varepsilon)L(\gamma) + \varepsilon L(\eta) < L(\gamma),$

in contradiction with the minimality of γ .

3. Lipschitz extensions

In this section, C and D will be two sets such that

D is a nonempty compact subset of \mathbb{R}^n , and

(4)
$$C$$
 is a closed subset of D containing ∂D ,

so that $D \setminus C$ is an open (possibly empty) set.

For a given continuous function $\psi \colon C \to \mathbb{R}$, we want to discuss the existence of a locally Lipschitz extension of ψ in D with the gradient constrained to belong to a convex set K.

It is well known that, in the case $D = \overline{\Omega}$, Ω bounded open subset of \mathbb{R}^n , $C = \partial \Omega$, and K satisfying (2), a necessary and sufficient condition for the existence of such an extension is

$$\psi(x) - \psi(y) \le L(\gamma) \qquad \forall x, y \in \partial\Omega, \ \forall \gamma \in \Gamma_{y,x}^{\Omega}$$

where $\Gamma_{y,x}^{\Omega}$ is the family of absolutely continuous curves $\gamma \colon [0,1] \to \overline{\Omega}$ such that $\gamma(t) \in \Omega$ for every $t \in (0,1), \ \gamma(0) = y$ and $\gamma(1) = x$ (see [26, Sect. 5] and [3]). Moreover, if Ω has a Lipschitz boundary, one can relax the above necessary and sufficient condition by taking paths in $\Gamma_{y,x}^{\overline{\Omega}}$ (see [26]). In order to handle more general situations, we will consider a continuous

function $\psi \colon C \to \mathbb{R}$ satisfying

(5)
$$\psi(x) - \psi(y) \le L(\gamma) \qquad \forall x, y \in C, \ \forall \gamma \in \Gamma_{y,x}^D,$$

where $\Gamma^{D}_{x_1,x_2}$ is the family of curves introduced in Definition 2.3.

This condition can be rephrased in the setting of the length space (D, d_L) , where $d_L: D \times D \to [0, +\infty]$ is the (possibly asymmetric) distance defined by

(6)
$$d_L(x,y) := \inf\{L(\gamma): \ \gamma \in \Gamma_{y,x}^D\}, \qquad x,y \in D.$$

(We refer to [9, Chap. 2] for the basic properties of length spaces; see in particular Remark 2.2.6 for the extension to possibly asymmetric metrics.) Namely, (5) is equivalent to requiring that ψ is a 1-Lipschitz function with respect to the d_L .

In the following lemma we recall the well known fact that (5) turns out to be a sufficient condition for the existence of minimal and maximal 1-Lipschitz extensions (w.r.t. d_L) of ψ . Notice that (5) is no more a necessary condition, as it can be seen when D is a closed segment in \mathbb{R}^2 .

Lemma 3.1. Let C, D satisfy (4), let K satisfy (2), and let $\psi: C \to \mathbb{R}$ be a continuous function satisfying (5). Then the functions

(7)
$$u^{\psi}(x) := \inf \left\{ \psi(y) + L(\gamma) : \ y \in C, \ \gamma \in \Gamma_{y,x}^{D} \right\}, \quad x \in D, \\ w^{\psi}(x) := \sup \left\{ \psi(y) - L(\gamma) : \ y \in C, \ \gamma \in \Gamma_{x,y}^{D} \right\}, \quad x \in D,$$

are continuous in D, and $u^{\psi} = w^{\psi} = \psi$ in C. Moreover, $u^{\psi}, w^{\psi} \in$ $W^{1,\infty}(\operatorname{int} D)$ and $Du^{\psi}, Dw^{\psi} \in K$ a.e. in $\operatorname{int} D$.

As a consequence of Lemma 3.1, the functions u^{ψ} and w^{ψ} belong to the functional space

$$Z_{\psi}:=\left\{u\in C(D)\cap W^{1,\infty}(\operatorname{int}D):\ u=\psi\ \text{on}\ C,\ Du\in K\ \text{a.e.\ in}\ \operatorname{int}D\right\}.$$

It should be noted that in the definition of $u^{\psi}(x)$ the infimum is taken over all paths in D joining $y \in C$ to x, while the supremum in the definition of w^{ψ} is taken over all paths in D joining x to $y \in C$. This asymmetry is needed in order to construct exactly the maximal and the minimal elements of Z_{ψ} , as it is stated in Lemma 3.3 below (see [6, Prop. 2.1]).

Definition 3.2. Given $x \in D$, we shall say that $y \in C$ is a visible point for x if the segment $[\![y,x[\![$]] is entirely contained in the open set $D \setminus C$. We shall denote by V_x^C the set of all visible points for x, that is $V_x^C = \{y \in C : [\![y,x[\![}]] \subset D \setminus C\}$.

Lemma 3.3. The functions u^{ψ} , w^{ψ} defined in (7) are characterized in the following way:

(8)
$$u^{\psi}(x) = \min\{\psi(y) + \rho^{0}(x - y) : y \in V_{x}^{C}\}, \qquad x \in D, \\ w^{\psi}(x) = \max\{\psi(y) - \rho^{0}(y - x) : y \in V_{x}^{C}\}, \qquad x \in D.$$

Moreover $w^{\psi} \leq u \leq u^{\psi}$ for every $u \in Z_{\psi}$.

Corollary 3.4. Let $x \in D \setminus C$ be a point of differentiability for u^{ψ} (resp. w^{ψ}). Then $Du^{\psi}(x)$ (resp. $Dw^{\psi}(x)$) belongs to ∂K .

Proof. Given $x \in D \setminus C$, let $y \in V_x^C$ be such that

$$u^{\psi}(x) = \psi(y) + \rho^{0}(x - y).$$

Setting $\xi = (x - y)/\rho^0(x - y)$, the fact that $u^{\psi} \in Z_{\psi}$ and Lemma 2.2 imply that

$$u^{\psi}(x-s\xi) - u^{\psi}(y) \le \rho^{0}(x-y) - s \,, \ u^{\psi}(x) - u^{\psi}(x-s\xi) \le s \quad \forall s \in [0,\rho^{0}(x-y)] \,.$$

Then we have

$$-s \le u^{\psi}(x - s\xi) - u^{\psi}(x) = u^{\psi}(x - s\xi) - u^{\psi}(y) - \rho^{0}(x - y) \le -s,$$

and hence

(9)
$$u^{\psi}(x - s\xi) = u^{\psi}(x) - s, \quad \forall s \in [0, \rho^{0}(x - y)].$$

If u^{ψ} is differentiable at x, (9) implies that $\langle Du^{\psi}(x), \xi \rangle = 1$, so that

$$\rho(Du^{\psi}(x)) = \max_{\eta \in K^0} \langle Du^{\psi}(x), \, \eta \rangle \ge 1.$$

On the other hand, by Lemma 3.1, $Du^{\psi}(x) \in K$, hence $Du^{\psi}(x)$ belongs to ∂K .

The proof for w^{ψ} is omitted, since it is entirely similar.

4. The existence result

Our main goal is to provide a full analysis of the Monge-Kantorovich system of PDEs

(10)
$$\begin{cases} -\operatorname{div}(v D\rho(Du)) = f & \text{in } \Omega, \\ \rho(Du) \le 1, \ v \ge 0 & \text{a.e. in } \Omega, \\ (1 - \rho(Du))v = 0 & \text{a.e. in } \Omega, \\ u = \phi & \text{pointwise on } \partial\Omega, \end{cases}$$

under the following assumptions:

- (H1) Ω is a nonempty bounded connected open subset of \mathbb{R}^n , $n \geq 1$.
- (H2) $f \in L^1(\Omega), f \geq 0$ a.e. in Ω .
- (H3) $\rho: \mathbb{R}^n \to \mathbb{R}$ is the gauge function of a set $K \subset \mathbb{R}^n$ satisfying (2).
- (H4) $\phi: \partial\Omega \to \mathbb{R}$ is a continuous function, satisfying the condition

$$\phi(x) - \phi(y) \le L(\gamma) := \int_0^1 \rho^0(\gamma'(t)) dt \qquad \forall x, y \in \partial\Omega, \ \forall \gamma \in \Gamma_{y,x},$$

where

$$\Gamma_{y,x} := \Gamma_{y,x}^{\overline{\Omega}} = \left\{ \gamma \in AC([0,1], \overline{\Omega}) : \ \gamma(0) = y, \ \gamma(1) = x \right\}.$$

Remark 4.1. In the first equation in (10) we understand $D\rho$ defined also at the origin since, by the third condition, v = 0 a.e. where Du = 0.

The functional setting for the unknowns (u, v) in (10) is $X_{\phi} \times L^{1}_{+}(\Omega)$, where

(11)
$$X_{\phi} := \{ u \in C(\overline{\Omega}) \cap W^{1,\infty}(\Omega) : u = \phi \text{ on } \partial\Omega, \ Du \in K \text{ a.e. in } \Omega \}$$
 and

$$L^1_+(\Omega):=\left\{v\in L^1(\Omega):\ v\geq 0\ \text{a.e. in }\Omega\right\}.$$

By a solution of (10) we mean a pair $(u,v) \in X_{\phi} \times L^{1}_{+}(\Omega)$ satisfying

(12)
$$\int_{\Omega} v \langle D\rho(Du), D\varphi \rangle = \int_{\Omega} f\varphi \qquad \forall \varphi \in C_0^{\infty}(\Omega)$$

and $(1 - \rho(Du))v = 0$ a.e. in Ω .

Let us define the Lax-Hopf function $u_{\phi} \colon \overline{\Omega} \to \mathbb{R}$ by

(13)
$$u_{\phi}(x) := \inf \left\{ \phi(y) + L(\gamma) : y \in \partial \Omega, \ \gamma \in \Gamma_{y,x} \right\}, \quad x \in \overline{\Omega},$$

which corresponds to the function u^{ψ} defined in (8) in the special case $D = \overline{\Omega}$, $C = \partial \Omega$, and $\psi = \phi$ (observe that $C = \partial \Omega \supseteq \partial \overline{\Omega} = \partial D$ and $\Omega \subseteq \operatorname{int} D$).

As a consequence of Lemmas 3.1, 3.3, and Corollary 3.4, we obtain the following result, which is well-known when Ω is a Lipschitz domain (see e.g. [26, Chap. 5]).

Theorem 4.2. If (H1), (H3) and (H4) are fulfilled, then $u_{\phi} \in X_{\phi}$ and $Du_{\phi} \in \partial K$ a.e. in Ω . Moreover, u_{ϕ} is the maximal element in X_{ϕ} , and it is characterized by

(14)
$$u_{\phi}(x) = \min\{\phi(y) + \rho^{0}(x - y) : y \in V_{x}\}, \qquad x \in \overline{\Omega},$$

where $V_x = \{ y \in \partial\Omega : \]\![y,x[\ \subset\Omega] \}.$

If $\phi = 0$, the Lax-Hopf function u_0 is the distance from $\partial\Omega$ with respect to the convex metric associated to ρ^0 , that is $u_0(x) = \min\{\rho^0(x-y) : y \in \partial\Omega\}$, $x \in \Omega$ (see [16] for a complete treatment of this subject).

More generally, if $\phi \colon \partial \Omega \to \mathbb{R}$ satisfies

(15)
$$\phi(x) - \phi(y) \le \rho^0(x - y) \qquad \forall x, y \in \partial\Omega,$$

the Lax–Hopf function u_{ϕ} coincides with

(16)
$$\tilde{u}(x) := \min\{\phi(y) + \rho^{0}(x - y) : y \in \partial\Omega\}, \qquad x \in \overline{\Omega}.$$

Namely, from (15) it is easy to check that $\tilde{u} = \phi = u_{\phi}$ on $\partial\Omega$. Moreover, from the representation formula (14) it is clear that $\tilde{u} \leq u_{\phi}$. On the other hand, let $x \in \Omega$ and let $y \in \partial\Omega$ be such that $\tilde{u}(x) = \phi(y) + \rho^{0}(x - y)$. Let $z \in [\![y,x]\!] \cap \partial\Omega$ be such that $[\![z,x]\!] \subset \Omega$. Then, by the representation formula (14) and by (15) we have that

$$\tilde{u}(x) \le u_{\phi}(x) \le \phi(z) + \rho^{0}(x-z) \le \phi(y) + \rho^{0}(z-y) + \rho^{0}(x-z)$$

= $\phi(y) + \rho^{0}(x-y) = \tilde{u}(x)$,

so that $\tilde{u}(x) = u_{\phi}(x)$.

The following example shows that if we merely require (H4), then in general $u_{\phi} \neq \tilde{u}$.

Example 4.3. Given $\epsilon \in (0,\pi)$ let us define the set

$$\Omega := \{ (\rho \cos \theta, \rho \sin \theta) \in \mathbb{R}^2 : 1 < \rho < 2, -\pi < \theta < \pi - \epsilon \},$$

and let $\phi \colon \partial \Omega \to \mathbb{R}$ be defined by $\phi = \theta$. If we choose $K = \overline{B}_1(0)$, then $\rho^0(\xi) = |\xi|$ for every $\xi \in \mathbb{R}^2$, and $L(\gamma)$ is the usual (Euclidean) length of a curve γ . A straightforward computation shows that ϕ satisfies (H4). On the other hand, it is also clear that ϕ does not satisfy (15). Namely, if we consider the two boundary points $y_1 = (-1, 0), y_2 = (\cos(\pi - \epsilon), \sin(\pi - \epsilon))$, then we have

$$\phi(y_2) - \phi(y_1) = 2\pi - \epsilon > 2\sin(\epsilon/2) = |y_2 - y_1|.$$

In this case, the function \tilde{u} defined in (16) does not belong to X_{ϕ} , since $\tilde{u}(y_2) \leq \phi(y_1) + |y_1 - y_2| < \phi(y_2)$ and hence it does not achieve the boundary value at y_2 .

We conclude this section with the following result due to S. Bianchini (see [6] and [7] for its generalization to the case of nonsmooth constraint sets K).

Theorem 4.4. There exists a weak solution $v_f \in L^1_+(\Omega)$ to the transport equation

(17)
$$-\operatorname{div}(v_f D\rho(Du_\phi)) = f \quad \text{in } \Omega.$$

As a corollary of Theorems 4.2 and 4.4 we conclude that (u_{ϕ}, v_f) is a solution to (10).

5. The minimum problem

In this section we shall discuss the connection between the system of PDEs (10) and the variational problem

(18)
$$\min \left\{ -\int_{\Omega} f u \, dx : u \in X_{\phi} \right\}.$$

It is clear that, since f is non-negative and u_{ϕ} is the maximal element in X_{ϕ} , then u_{ϕ} is a solution to (18). We shall show that (10) is the Euler-Lagrange necessary and sufficient condition for the minimum problem (18). This property, which is the analogous to the duality in Mass Transport Theory, will be useful in the proofs of the uniqueness of the solutions of (10).

As a preliminary step, we need to enrich the class of test functions allowed in (12).

Lemma 5.1. Let $u \in C(\overline{\Omega}) \cap W^{1,\infty}(\Omega)$, u = 0 on $\partial\Omega$. Then there exists a sequence $(\psi_j) \subset C_0^{\infty}(\Omega)$ such that $\psi_j \to u$ and $D\psi_j \to Du$ a.e. in Ω , and $\|\psi_j\|_{\infty} \leq \|u\|_{\infty}$, $\|D\psi_j\|_{\infty} \leq 3\|Du\|_{\infty}$ for every $j \in \mathbb{N}$.

Proof. We use a standard truncation argument. Let $G \in C^1(\mathbb{R})$ such that $0 \leq G(t) \leq |t|$ for every $t \in \mathbb{R}$, G(t) = 0 for $|t| \leq 1$, G(t) = |t| for $|t| \geq 2$, $|G'(t)| \leq 3$ for every $t \in \mathbb{R}$.

Let us define $u_j := G(ju)/j$, so that $u_j \to u$ pointwise in $\overline{\Omega}$. Moreover, $Du_j = G'(ju)Du$, hence $Du_j \to Du$ a.e. in Ω and $|Du_j| \le 3|Du|$.

Let φ_{ϵ} be the standard family of mollifiers in \mathbb{R}^n . Since u_j has compact support in Ω , we can choose a sequence $\varepsilon_j \searrow 0$, such that the sequence $\psi_j := \varphi_{\epsilon_j} * u_j$ has the required properties.

Corollary 5.2. $C_0^{\infty}(\Omega)$ is dense in the space $X := \{u - w : u, w \in X_{\phi}\}$ with respect to the weak* topology of $W^{1,\infty}(\Omega)$. As a consequence, if $(u,v) \in X_{\phi} \times L^1_+(\Omega)$ satisfies (12), then

(19)
$$\int_{\Omega} v \langle D\rho(Du), Dw - Du \rangle = \int_{\Omega} f(w - u) \qquad \forall w \in X_{\phi}.$$

We are now in a position to prove that the Monge-Kantorovich system (10) is the Euler-Lagrange condition for the minimum problem (18).

Theorem 5.3. The minimum problem (18) and the system of PDEs (10) are equivalent in the following sense.

- (i) $u \in X_{\phi}$ is a solution to (18) if and only if there exists $v \in L^1_+(\Omega)$ such that (u, v) is a solution to (10).
- (ii) Let $u \in X_{\phi}$ be a solution to (18). Then (u, v) is a solution to (10) if and only if (u_{ϕ}, v) is a solution to (10).

Proof. Let us denote by I_K the indicator function of the set K, that is

$$I_K(x) := \begin{cases} 0 & \text{if } x \in K, \\ +\infty & \text{if } x \in \mathbb{R}^n \setminus K, \end{cases}$$

so that

$$F(u) := \int_{\Omega} [I_K(Du) - fu] dx = -\int_{\Omega} fu dx, \quad \forall u \in X_{\phi}.$$

Since ρ is differentiable in $\mathbb{R}^n \setminus \{0\}$, the subgradient of the indicator function can be explicitly computed, obtaining

(20)
$$\partial I_K(\xi) = \begin{cases} \{\alpha D \rho(\xi) \colon \alpha \ge 0\} & \text{if } \xi \in \partial K, \\ \emptyset & \text{if } \xi \notin K, \\ \{0\} & \text{if } \xi \in \text{int } K \end{cases}$$

(see e.g. [29, Sect. 23]).

Hence, if (u, v) is a solution to (10), then $v(x)D\rho(Du(x)) \in \partial I_K(Du(x))$ for a.e. $x \in \Omega$, so that, for every $w \in X_{\phi}$

(21)
$$F(w) - F(u) \ge \int_{\Omega} v \langle D\rho(Du), Dw - Du \rangle dx - \int_{\Omega} f(w - u) dx = 0$$
,

where the last equality follows from Corollary 5.2. This proves that u is a solution to (18).

Assume now that $u \in X_{\phi}$ is a minimizer for F, so that $f(u - u_{\phi}) = 0$ a.e. in Ω , due to the maximality of u_{ϕ} in X_{ϕ} and the fact that $f \geq 0$. By Corollary 5.2 we can choose $u - u_{\phi}$ as test function in the transport equation (17) solved by u_{ϕ} , getting

$$0 = \int_{\Omega} v_f \langle D\rho(Du_\phi), Du - Du_\phi \rangle dx = -\int_{\Omega} v_f (1 - \langle D\rho(Du_\phi), Du \rangle) dx.$$

On the other hand, by Theorem 2.1(ii) and the fact that $Du \in K$ a.e. in Ω , we have

$$1 - \langle D\rho(Du_{\phi}(x)), Du(x) \rangle \ge 0$$
, a.e. $x \in \Omega$,
 $1 - \langle D\rho(Du_{\phi}(x)), Du(x) \rangle = 0 \iff D\rho(Du_{\phi}(x)) = D\rho(Du(x))$,

so that $v_f D\rho(Du_\phi) = v_f D\rho(Du)$ and $v_f(1-\rho(Du)) = 0$ a.e. in Ω , that is (u, v_f) is a solution of (10). This concludes the proof of (i).

Let us prove (ii). The previous computation shows that if (u_{ϕ}, v) is a solution of (10), and $u \in X_{\phi}$ is a solution of the minimum problem (18), then also (u, v) is a solution of (10). Finally, let (u, v) be a solution to (10). Upon observing that $v\rho(Du) = v$ a.e. in Ω , and choosing $u - u_{\phi}$ as test function in the weak formulation of the transport equation

$$-\operatorname{div}(v D\rho(Du)) = f,$$
 in Ω

we conclude that the opposite implication holds.

6. Uniqueness of the v-component

Thanks to Theorem 5.3(ii), the uniqueness of the v component of the solution of the system of PDEs (10) will follow from the fact that the function v_f introduced in Theorem 4.4 is the unique weak solution in $L^1_+(\Omega)$ of the equation

$$-\operatorname{div}(v \, D\rho(Du_{\phi})) = f.$$

We shall use an explicit representation of v_f proved in [6] and, in order to explain it, we need to introduce some additional tools, related to the directions where the function u_{ϕ} has the maximal slope.

Recalling the representation formula (14) for u_{ϕ} , for every $x \in \overline{\Omega}$ let us define the projections of x by

$$\Pi(x) := \{ y \in V_x \colon \ u_{\phi}(x) = \phi(y) + \rho^0(x - y) \}$$

and, for every $x \in \Omega$, let $\Delta(x)$ be the set of directions through x

$$\Delta(x) := \left\{ \frac{x - y}{\rho^0(x - y)} : y \in \Pi(x) \right\}, \qquad x \in \Omega.$$

Let $D \subset \Omega$ be the set of points with multiple projections, that is

(23)
$$D := \{ x \in \Omega : \ \Delta(x) \text{ is not a singleton} \},$$

and for every $x \in \Omega \setminus D$, let p(x) and d(x) denote the unique elements in $\Pi(x)$ and $\Delta(x)$ respectively. Finally, let $\tau(x)$ be defined by (24)

$$\tau(x) := \begin{cases} \sup\{t \ge 0; \ u_{\phi}(x + sd(x))) = u_{\phi}(x) + s \ \forall s \in [0, t]\}, & x \in \Omega \setminus D, \\ 0 & x \in D, \end{cases}$$

and let J be the set

(25)
$$J := \bigcup_{x \in \Omega} q(x), \quad q(x) := x + \tau(x)d(x),$$

where we understand that q(x) = x if $x \in D$.

Definition 6.1. We shall call transport ray through $x \in \Omega$ any segment $[p, q(x)], p \in \Pi(x)$.

It is clear from the definition that, if $x \in \Omega \setminus D$, then there is a unique transport ray [p(x), q(x)] through x. On the other hand, if $x \in D$, any segment [p, x], with $p \in \Pi(x)$, is a transport ray through x.

The transport rays correspond to the segments where u_{ϕ} grows linearly with maximal slope, as stated in the following lemma. Hence in the sandpiles model, they correspond to the directions where the matter is allowed to run down, that is where the v-component could be nonzero.

Lemma 6.2. For every $x \in \Omega$ and $y \in \Pi(x)$ we have

(26)
$$u_{\phi}(z) = \phi(y) + \rho^{0}(z - y) \qquad \forall z \in [\![y, x]\!].$$

Moreover, if $x \in \Omega \setminus D$, then

(27)
$$u_{\phi}(z) = \phi(p(x)) + \rho^{0}(z - p(x)) \qquad \forall z \in [\![p(x), q(x)]\!].$$

Proof. The identity (26) follows from (9), while (27) for $x \in \Omega \setminus D$ follows from (26) when $z \in [p(x), x]$, and from the very definition of q(x) when $z \in [x, q(x)]$.

Corollary 6.3. Given $x \in \Omega \setminus D$, let [p(x), q(x)] be the transport ray through x. Then $\Pi(z) = \{p(x)\}$ for every $z \in [p(x), q(x)]$.

Proof. From (27) it is clear that $p(x) \in \Pi(z)$ for every $z \in [p(x), q(x)]$. Assume by contradiction that there exists $z \in [p(x), q(x)]$ such that $\Pi(z)$ contains a point $y \neq p(x)$.

We claim that the three points y, z and q(x) cannot be aligned. Namely, we cannot have $z \in [y, q(x)[$, for otherwise we would have y = p(x). On the other hand, we cannot have $q(x) \in [y, z]$: indeed, in this case, by (27) we would have $u_{\phi}(q(x)) = \phi(y) + \rho^{0}(q(x) - y)$, while, by (26), we know that $u_{\phi}(q(x)) = u_{\phi}(z) + \rho^{0}(q(x) - z)$. Putting together these information, we get

$$u_{\phi}(z) = \phi(y) + \rho^{0}(z - y) = \phi(y) - \rho^{0}(z - q(x)) + \rho^{0}(q(x) - y)$$
$$= u_{\phi}(z) + \rho^{0}(q(x) - z) + \rho^{0}(z - q(x)) > u_{\phi}(z),$$

a contradiction.

Since the curve $[y, z] \cup [z, q(x)]$, joining y to q(x) and lying in Ω except for its endpoints, is not a segment, by Lemma 2.4 it cannot be a geodesic, so that

$$u_{\phi}(q(x)) < \phi(y) + \rho^{0}(z-y) + \rho^{0}(q(x)-z).$$

Taking into account (27) we get

$$u_{\phi}(z) = u_{\phi}(q(x)) - \rho^{0}(q(x) - z) < \phi(y) + \rho^{0}(z - y),$$

a contradiction with $y \in \Pi(z)$.

At this stage we are dealing with three meaningful sets of singular points associated to u_{ϕ} :

- Σ , the sets of points in Ω where u_{ϕ} is not differentiable;
- D, the sets of points in Ω with multiple projections;
- J, the set of endpoints in Ω of the transport rays.

The set J plays a central role in optimal transport theory. As a matter of fact, the disintegration of the Lebesgue measure in the reference set Ω along transport rays, can be achieved only if J has Lebesgue measure zero, and this is one of the main tools of the theory.

If Ω and K are sufficiently smooth (e.g. both of class C^2 , and K with strictly positive curvatures), one can prove that D coincides with Σ (see, e.g. [14, 16]), so that, by Rademacher's Theorem, D has zero Lebesgue measure. Moreover, using the methods of Differential Geometry, it is possible to prove that $J = \overline{\Sigma}$ and to characterize $J \setminus \Sigma$ in terms of optimal focal points. As a consequence one can prove that J has zero Lebesgue measure and give detailed rectifiability results (see, e.g., [16, 21]).

If Ω and K are not regular, the set $\overline{\Sigma}$ may have positive Lebesgue measure (see [27] for an example where $\Omega \in C^{1,\alpha}$, $\alpha < 1$, and K is the unit ball), and J can be a proper subset of $\overline{\Sigma}$.

We recall here some known properties of the singular sets, adding some remarks in order to complete the description of the relationships between these sets. The known results, proved in [6], are collected in the following proposition.

Proposition 6.4. The sets Σ , D and J satisfy the following properties.

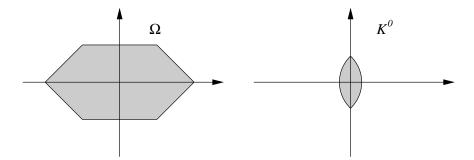


FIGURE 2. The sets Ω and K^0 of Example 6.5

- (i) $D \subseteq \Sigma$. More precisely, if u_{ϕ} is differentiable at $x \in \Omega$, then x has a unique projection and $\Delta(x) = \{D\rho(Du_{\phi}(x))\}.$
- (ii) The set J has zero Lebesque measure.
- (iii) D is a (n-1) rectifiable set.

Remark 6.5. We recall that $D = \Sigma$ if K is strictly convex. Proposition 6.4(i) states that $D \subseteq \Sigma$, in general. The following example shows that it can happen that $D \neq \Sigma$. Let us consider the set

$$K^{0} = \left\{ (x, y) \in \mathbb{R}^{2} \colon \left(|x| + \frac{\sqrt{2}}{2} \right)^{2} + y^{2} \le 1 \right\}$$

and let K be its polar set. The set K^0 is strictly convex but it is not of class C^1 , hence the set K is of class C^1 but it is not strictly convex (see Theorem 2.1). Let Ω be the hexagon

$$\Omega = \{(x, y) \in \mathbb{R}^2 \colon |y| < 2 - |x|, \ |y| < 1\}$$

(see Figure 2). The Lax-Hopf function u_0 associated to K and to the boundary datum $\phi = 0$ is the distance function to the boundary of Ω with respect to the convex metric ρ^0 , which can be explicitly computed as

$$u_0(x) = \begin{cases} \sqrt{2}(1 - |y|), & \text{if } (x, y) \in \Omega \cap \{|x| \le 1\}, \\ \sqrt{2}(2 - |x| - |y|), & \text{if } (x, y) \in \Omega \cap \{|x| > 1\}, \end{cases}$$

so that $\Sigma \setminus D = \{(x, y) \in \Omega \colon |x| = 1, y \neq 0\}.$

Remark 6.6. Since D is a subset of J, then by Proposition 6.4(ii) it has Lebesgue measure zero. Notice that D can be a proper subset of J even in the regular case. For example, let Ω be the ellipsis $\{(x,y) \in \mathbb{R}^2; x^2/a^2 + y^2/b^2 < 1\}$, where 0 < b < a, and let $K = \overline{B}_1(0)$. The points $P = ((a^2 - b^2)/a, 0)$ and $Q = (-(a^2 - b^2)/a, 0)$ are the centers of curvature of $\partial \Omega$ at (a,0) and (-a,0) respectively. It can be checked that $J = [\![Q,P]\!]$ whereas $D = [\![Q,P]\!]$.

Moreover, since J need not be closed (see [6, Remark 5.2]), $J \neq \overline{D}$ in general. The following proposition shows that $J \subseteq \overline{D}$.

Proposition 6.7. For every $x \in \Omega \setminus \overline{D}$ we have that $\tau(x) > 0$, where τ is the function defined in (24).

Proof. Let us fix $x \in \Omega \setminus \overline{D}$. We can assume without loss of generality that x = 0, and $d(0) = -\lambda e_n$, $\lambda > 0$. Since the map $z \mapsto d(z)$ is continuous in the open set $\Omega \setminus \overline{D}$ (see [6, Prop. 3.2]), there exists r > 0 such that $\overline{B}_r(0) \subseteq \Omega \setminus \overline{D}$, and $\langle d(z), e_n \rangle < -\lambda/2$ for every $z \in \overline{B}_r(0)$.

Given $\delta \in (0, r/2)$, let us define the function $\psi \colon \overline{B}_r(0) \to \mathbb{R}^n$ by

(28)
$$\psi(z) := z + \left(\frac{2\delta - \langle z, e_n \rangle}{\langle d(z), e_n \rangle}\right) d(z),$$

so that $\langle \psi(z), e_n \rangle = 2\delta$. Hence ψ maps $\overline{B}_r(0)$ into the hyperplane $H := \{x \in \mathbb{R}^n : \langle x, e_n \rangle = 2\delta\}$. The continuity of d(z) in $\overline{B}_r(0)$ ensures that we can choose δ in such a way that $\psi(z) \in \overline{B}_r(0)$ for every $z \in \overline{B}_\delta(0)$.

Summarizing, $\psi \colon \overline{B}_r(0) \to H$ is a continuous map, $\psi(\overline{B}_{\delta}(0)) \subset \overline{B}_r(0)$, and $\psi(z)$ belongs to the ray $\{z - td(z) : t \geq 0\}$. Moreover, for every $z \in \overline{B}_{\delta}(0)$, we have that $\llbracket \psi(z), z \rrbracket \subset \overline{B}_r(0) \subset \Omega$, hence $\psi(z) \in \llbracket p(z), z \llbracket$ and, by Corollary 6.3, we infer that $d(z) = d(\psi(z))$.

Since ψ maps $\partial B_{\delta}(0)$ into the (n-1)-dimensional hyperplane H, by the Borsuk–Ulam lemma (see [20, Cor. 4.2]) there exists $w \in \partial B_{\delta}(0)$ such that $\psi(w) = \psi(-w)$. Furthermore $d(w) = d(\psi(w)) = d(\psi(-w)) = d(-w)$, so that the equality $\psi(w) = \psi(-w)$ implies that

$$w = \frac{\langle w, e_n \rangle}{\langle d(w), e_n \rangle} d(w),$$

i.e., w is parallel to d(w). But the transport ray through w and -w contains the origin, hence d(w) = d(-w) = d(0) and $\tau(0) \ge \delta$.

For the reader's convenience we collect here the results concerning the relationships between the singular sets.

Corollary 6.8. The sets Σ , D and J have zero Lebesgue measure. Moreover, $D \subset \Sigma$ and $D \subset J \subset \overline{D}$, with possibly strict inclusions.

We are now in a position to briefly recall the explicit representation of the function v_f appearing in (17), one of the main results in [6].

Theorem 6.9. There exists a positive function $\alpha \colon \Omega \to (0, +\infty)$, absolutely continuous along almost every transport ray, such that an explicit solution $v_f \in L^1_+(\Omega)$ of (17) is given by

(29)
$$v_f(x) = \int_0^{\tau(x)} f(x + td(x)) \frac{\alpha(x + td(x))}{\alpha(x)} dt, \quad a.e. \ x \in \Omega.$$

More precisely, for a.e. $x \in \Omega \setminus D$ the function v_f is locally absolutely continuous along the ray $t \mapsto x + td(x)$, $t \in [0, \tau(x))$, and satisfies

(30)
$$\begin{cases} \frac{d}{dt} \left[v_f(x + td(x)) \alpha(x + td(x)) \right] = -f(x + td(x)) \alpha(x + td(x)), \\ \lim_{t \to \tau(x)^-} v_f(x + td(x)) \alpha(x + td(x)) = 0. \end{cases}$$

Proof. See [6], Theorems 5.7 and 7.1.

Remark 6.10. The reader may want to compare the representation formula (29) with the one proved in the regular homogeneous case (see [10, 11, 16]), where $\alpha(x)$ can be explicitly written in terms of the distance from the boundary $u_0(x)$ and of the anisotropic curvatures $\kappa_1, \ldots, \kappa_{n-1}$ of $\partial\Omega$ at p(x). More precisely, one has

$$\alpha(x) = \prod_{j=1}^{n-1} \left[1 - u_0(x) \, \kappa_j(p(x)) \right], \qquad x \in \Omega \setminus D.$$

We shall find a necessary and sufficient condition for having that v_f is the unique solution to (22). The following example shows why the uniqueness of the v-component may fail.

Example 6.11. Let $\Omega = (0,1) \times (0,1) \subseteq \mathbb{R}^2$, $K = \overline{B}_1(0)$, and let $\phi(x,y) = y$ on $\partial\Omega$, so that $u_{\phi}(x,y) = y$ in Ω . Then every function of the form v(x,y) = 1 - y + c(x), $c(x) \geq 0$ for $x \in [0,1]$, is a non-negative solution of the equation $-\operatorname{div}(v D u_{\phi}) = 1$. Notice that in this case $v_f(x,y) = 1 - y$ in Ω , so that (22) has infinitely many solutions of the form $v = v_f + \tilde{v}$, with $\operatorname{div}(\tilde{v} D u_{\phi}) = 0$. The existence of non-trivial variations of v_f is due to the fact that Ω is covered by transport rays with both endpoints on $\partial\Omega$.

Now it should be clear that, in order to discuss the uniqueness of the v-component of the solutions to (22), we need to introduce the set

$$T := \bigcup_{(p,q) \in E}]\![p,q[\![\,,$$

where the set E is defined by

$$E := \{ (p,q) \in \partial\Omega \times \partial\Omega; \ p \neq q, \ \|p,q\| \subseteq \Omega, \ \phi(q) = \phi(p) + \rho^0(q-p) \}.$$

In Example 6.11, $E = \{((x,0),(x,1)), x \in (0,1)\}$ and $T = \Omega$. In general T is the union of the transport rays ending on $\partial\Omega$. On this set every function in X_{ϕ} is forced to have maximal slope and hence to coincide with u_{ϕ} .

Lemma 6.12. Let $(p,q) \in E$ be given. Then the following hold.

- i) $\Pi(x) = \{p\}$ and q(x) = q for every $x \in [p, q]$.
- ii) If $u \in X_{\phi}$, then $u(x) = \phi(p) + \rho^{0}(x p)$ for every $x \in [p, q]$. As a consequence every $u \in X_{\phi}$ coincides with u_{ϕ} on the set T.

Proof. Let us start by proving that $p \in \Pi(x)$ for every $x \in [p, q]$. Namely

$$u_{\phi}(x) \le \phi(p) + \rho^{0}(x - p) = \phi(q) - \rho^{0}(q - p) + \rho^{0}(x - p)$$

= $\phi(q) - \rho^{0}(q - x) \le u_{\phi}(x)$,

where the last inequality follows from Lemma 2.2. Hence $u_{\phi}(x) = \phi(p) + \rho^{0}(x-p)$, so that $p \in \Pi(x)$.

Actually we can prove that $\Pi(x) = \{p\}$ for every $x \in [p, q]$ using the same argument of the proof of Corollary 6.3: if there exists $p_1 \neq p$, $p_1 \in \Pi(x)$,

then the three points p_1 , x and q cannot be aligned, so that

$$\phi(q) - \phi(p_1) < \rho^0(q - x) + \rho^0(x - p_1)$$

and

$$u_{\phi}(x) = \phi(p_1) + \rho^0(x - p_1) > \phi(q) - \rho^0(q - x) = \phi(p) + \rho^0(x - p) = u_{\phi}(x),$$

a contradiction.

Since $q \in \partial \Omega$ and, for every $x \in [p], q[n]$, $\Pi(x) = \{p\}$, it is clear that q(x) = q.

Finally, let $u \in X_{\phi}$ and $x \in [p,q]$ be given. Then, using Lemma 2.2 we get

$$\phi(p)+\rho^0(x-p)=\phi(q)-\rho^0(q-x)\leq u(x)\leq \phi(p)+\rho^0(x-p)\,,$$
 and (ii) follows.

Theorem 6.13. The function v_f defined in (29) is the unique solution of (22) if and only if the set T has zero Lebesgue measure.

Proof. The uniqueness of v_f when T has zero Lebesgue measure is proved in [6, Prop. 7.3]. The proof is based on the fact that every weak solution $v \in L^1_+(\Omega)$ to (22) is locally absolutely continuous along almost every transport ray, and satisfies

$$\frac{d}{dt} \left[v(x + td(x)) \alpha(x + td(x)) \right] = -f(x + td(x)) \alpha(x + td(x)).$$

Moreover, if the endpoint q(x) of the transport ray through x belongs to Ω , the $v \cdot \alpha$ satisfies the initial condition

$$\lim_{t \to \tau(x)^{-}} v_f(x + td(x)) \alpha(x + td(x)) = 0$$

so that, by Theorem 6.9, $v = v_f$ along that ray.

Let us now assume that T has positive Lebesgue measure. It is clear that the function

$$w_{\phi}(x) := \min\{-\phi(y) + \rho^{0}(y - x) : y \in V_{x}\}, \qquad x \in \overline{\Omega}$$

is the Lax-Hopf function corresponding to the geometry induced by the convex set -K and the boundary datum $-\phi$. Assume that u_{ϕ} and w_{ϕ} are differentiable at a point $x \in [p, q[$, with $(p, q) \in E$. From Proposition 6.4 we have that $Du_{\phi}(x) = (q-p)/\rho^0(q-p)$. The same argument, used with $-\phi$ as boundary datum and the geometry induced by -K, shows that

$$D\rho_{-}(Dw_{\phi}(x)) = \frac{p-q}{\rho_{-}^{0}(p-q)} = -\frac{q-p}{\rho^{0}(q-p)} = D\rho(Du_{\phi}(x)),$$

where ρ_{-} and ρ_{-}^{0} denote respectively the gauge functions of the convex sets -K and $-K^{0}$.

Let us consider the indicator function of the set T,

$$g(x) := \begin{cases} 1, & x \in T, \\ 0, & x \in \Omega \setminus T. \end{cases}$$

By Theorem 4.4, there exist two nonnegative functions v_q^+, v_q^- such that

$$-\operatorname{div}(v_g^+ D\rho(Du_\phi)) = g, \quad -\operatorname{div}(v_g^- D\rho_-(Dw_\phi)) = g, \quad \text{in } \Omega.$$

Looking at the explicit representations of v_g^+ and v_g^- as in (29), we easily

$$v_g^+=v_g^-=0 \text{ a.e. in } \Omega\setminus T\,, \qquad v_g^+>0, \ v_g^->0 \text{ a.e. in } T\,,$$
 so that, for every $\varphi\in C_0^\infty(\Omega)$, we obtain

$$\int_{\Omega} (v_g^+ + v_g^-) \langle D\rho(Du_\phi), D\varphi \rangle dx = \int_{T} (v_g^+ + v_g^-) \langle D\rho(Du_\phi), D\varphi \rangle dx$$

$$= \int_{T} v_g^+ \langle D\rho(Du_\phi), D\varphi \rangle dx - \int_{T} v_g^- \langle D\rho_-(Dw_\phi), D\varphi \rangle dx$$

$$= \int_{\Omega} g \varphi dx - \int_{\Omega} g \varphi dx = 0.$$

Hence for every $\lambda \geq 0$ the function $v_{\lambda} = v_f + \lambda(v_g^+ + v_g^-)$ is a non-negative weak solution to (22).

Remark 6.14. It can be easily checked that if the datum ϕ satisfies (H4) with strict inequality holding for every $x \neq y$, then the Lebesgue measure of T is zero, so that v_f is the unique v-component allowed in (10).

Remark 6.15. If the Lebesgue measure of $\Omega \setminus T$ is zero, then by Lemma 6.12 we have that $X_{\phi} = \{u_{\phi}\}$. On the other hand, it can be easily checked that the converse implication is also true. In [5] (see also [6]) it is proved that either $X_{\phi} \neq \{u_{\phi}\}$ or there exists a weak solution $v \in L^{1}_{loc}(\Omega)$ to

$$\operatorname{div}(v D u_{\phi}) = 0$$

such that v > 0 a.e. in Ω . In the proof of Theorem 6.13 we have shown that, as soon as T has positive Lebesgue measure, it is possible to construct a solution $v \in L^1_+(\Omega)$ to (31) such that v > 0 a.e. on T.

7. Uniqueness of the u-component

In Theorem 5.3 we have shown that the set of admissible u-components of the solutions to (10) coincides with the set of solutions to the minimum problem (18).

In this section we shall construct, using Lemma 3.1, the minimal solution to (18), and we shall give a necessary and sufficient condition in order to have that this function equals the maximal solution u_{ϕ} in Ω .

Let us define the function $u_f : \overline{\Omega} \to \mathbb{R}$ by

(32)
$$u_f(x) := \sup\{u_\phi(z) - L(\gamma) : z \in \partial\Omega \cup \operatorname{supp}(f), \ \gamma \in \Gamma_{x,z}\},\$$

where $\operatorname{supp}(f)$ is the complement in $\overline{\Omega}$ of the union of all relatively open subsets $A \subseteq \Omega$ such that f = 0 a.e. in A.

Proposition 7.1. The function u_f is characterized by

(33)
$$u_f(x) = \max\{u_\phi(z) - \rho^0(z - x) : z \in W_x\},\,$$

where $W_x := \{ y \in \partial\Omega \cup \operatorname{supp}(f) : \|y, x\| \subset \Omega \setminus \operatorname{supp}(f) \}, \text{ and satisfies}$

- (i) $u_f \in X_{\phi}$;
- (ii) $u_f(x) = u_\phi(x)$ for every $x \in \text{supp}(f) \cup T$;
- (iii) given $u \in X_{\phi}$, we have $u = u_{\phi}$ on supp(f) if and only if $u_f \leq u \leq u_{\phi}$ in Ω .

Proof. The function u_f is the function w^{ψ} defined in (7) corresponding to $D = \overline{\Omega}$, $C = \partial \Omega \cup \text{supp}(f)$, and $\psi = u_{\phi}$ on C. Clearly the function ψ satisfies (5) on C, since u_{ϕ} satisfies that condition on $\overline{\Omega}$. Then the stated properties follow from Lemmas 3.1, 3.3 and 6.12.

Theorem 7.2. A function $u \in X_{\phi}$ is a solution to (18) if and only if $u_f \leq u \leq u_{\phi}$. Moreover, the function u_f coincides with u_{ϕ} in Ω (and hence u_{ϕ} is the unique solution to (18)) if and only if $J \cap \Omega \subseteq \text{supp}(f)$.

Proof. The first assertion follows from Proposition 7.1(iii) and from the fact that $u \in X_{\phi}$ is a solution to (18) if and only if $u = u_{\phi}$ on $\operatorname{supp}(f)$. Let us prove the uniqueness result. Assume that $J \cap \Omega \subseteq \operatorname{supp}(f)$, so that $J \subset \partial\Omega \cup \operatorname{supp}(f)$ and, from Proposition 7.1(ii), $u_f = u_{\phi}$ on J. Let $x \in \Omega \setminus J$ be given, and let $q(x) \in J$ be the endpoint of the ray through x. We have

$$u_{\phi}(x) = u_{\phi}(q(x)) - \rho^{0}(q(x) - x) = u_{f}(q(x)) - \rho^{0}(q(x) - x) \le u_{f}(x) \le u_{\phi}(x)$$
, and hence $u_{f}(x) = u_{\phi}(x)$.

Assume now by contradiction that $u_f = u_\phi$ in Ω and that there exists $x_0 \in J \cap \Omega$ such that $x_0 \notin \text{supp}(f)$. Let $z \in W_{x_0}$ and $y \in V_{x_0}$ be such that

(34)
$$u_f(x_0) = u_{\phi}(z) - \rho^0(z - x_0),$$

(35)
$$u_{\phi}(x_0) = u_{\phi}(y) + \rho^0(x_0 - y).$$

Notice that $y \neq z$, otherwise we get $-\rho^0(y-x_0) = \rho^0(x_0-y)$. Hence x_0, y and z are three distinct points. Moreover we have that $x_0 \in [y, z[$, otherwise by Lemma 2.4 we should have

$$u_{\phi}(z) - u_{\phi}(y) < \rho^{0}(z - x_{0}) + \rho^{0}(x_{0} - y)$$

and

$$u_{\phi}(x_0) = u_{\phi}(z) - \rho^0(z - x_0) < u_{\phi}(y) + \rho^0(x_0 - y) = u_{\phi}(x_0),$$

a contradiction.

Finally, the fact that $x_0 \in [y, z]$ implies that

$$u_{\phi}(z) = u_{\phi}(x_0) + \rho^0(z - x_0) = u_{\phi}(y) + \rho^0(x_0 - y) + \rho^0(z - x_0)$$

= $u_{\phi}(y) + \rho^0(z - y)$,

hence [y, z[is a transport ray through x_0 , so that $x_0 \notin J$.

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